

# 1 On the measurement of stability parameter over complex 2 mountainous terrain

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9 **Abstract.** Atmospheric stability has a significant effect on wind shear and turbulence intensity, and these variables,  
10 in turn, have a direct impact on wind power production and loads on wind turbines. It is therefore important to know  
11 how to characterize atmospheric stability in order to make better energy yield estimation in a wind farm.

12 Based on research grade meteorological mast at Alaiz (CENER's Test Site in Navarre, Spain) named MP5, this work  
13 compares and evaluates different instrument set-ups and methodologies for stability characterization, namely: the  
14 Obukhov parameter, measured with a sonic anemometer, and the bulk Richardson number based on two temperature  
15 and one wind speed measurements. The methods are examined considering their theoretical background,  
16 implementation complexity, instrumentation requirements, and practical use in connection to wind energy  
17 applications. The sonic method provides a more precise local measurement of stability while the bulk Richardson is  
18 a simpler, robust and cost-effective technique to implement in wind assessment campaigns. Using the sonic method  
19 as a benchmark, it is shown that to obtain reliable bulk Richardson measurements in onshore sites it is necessary  
20 install one of the temperature sensors close to the ground where the temperature gradient is stronger.

## 21 1. Introduction

22 The vertical wind profile and the turbulence intensity in the atmospheric boundary layer (ABL) are two of  
23 the main physics aspects driving wind energy production and turbine loads. The vertical wind profile is especially  
24 important since rotors are getting bigger and hub heights are getting higher making it invaluable to know the wind  
25 speed at hub height. The vertical wind profile shape and turbulence intensity can directly influence wind turbine  
26 production but also wind turbine loads, affecting the wind turbines lifetime. Despite the fact that the IEC standard  
27 (IEC61400-1 (ED4) 2019, 2019) specifies a power law vertical model independent of atmospheric stability to  
28 perform load calculations, the dependence of this and, in turn, the turbulence intensity with atmospheric stability is  
29 widely demonstrated (Emeis, 2013; Lange et al., 2004b; Peña and Hahmann, 2012). In addition several studies  
30 have demonstrated the impact of atmospheric stability on wind resource assessment (Lange et al., 2004a), wind  
31 turbine power curves and Annual Energy Production (AEP) calculations (Martin et al., 2016; Schmidt et al.,  
32 2016); wind turbine loads (Kelly et al., 2014; Sathe et al., 2013) and wind turbine wakes (Abkar and Porté-Agel,  
33 2015; Hansen et al., 2010; Macheffaux et al., 2016). This is why the wind industry is developing models and  
34 methods to include the effect of atmospheric stability in the layout design and energy yield assessment. These  
35 methodologies and models require the characterization of the probability distribution of atmospheric stability at  
36 each site. Therefore different methods and parameters are used to describe atmospheric stability without an  
37 industry-wide convention about which one is the most appropriate.

38 According to Monin and Obukhov similarity theory (MOST) (Foken, 2006; Monin and Obukhov, 1954) stability  
39 can be estimated in terms of inverse of Obukhov length that can be calculated with vertical fluxes of heat and  
40 momentum obtained with the eddy covariance method. To obtain the necessary high-frequency measurements of  
41 wind speed vector components and temperature, sonic anemometers are used, which is why this calculation method  
42 is called "sonic method".

43 Another measure for stability is the Richardson number that as Bardal (Bardal et al., 2018) explains according to  
44 Stull book (Stull, 1989) has several formulations: the flux Richardson number, gradient Richardson number and  
45 bulk Richardson number. The latter is based on one height wind speed measurement and two temperature  
46 measurements, one from the air at one height and the other from the ground or water surface.

47 In the wind energy context some studies have been done about how to measure the stability and their influence in  
48 the turbulence intensity and vertical wind profile. However, most of these studies have been carried out in offshore  
49 sites (Peña and Hahmann, 2012; Sanz Rodrigo et al., 2015; Sathe et al., 2011) finding relationships (Grachev and  
50 Fairall, 1997) between the Obukhov length and the Richardson bulk number that, facilitate the characterization of  
51 stability without the need of sonic anemometer. This is convenient, because although the sonic anemometer has  
52 many advantages (Cuerva et al., 2006), they add complexity, in terms of use and data management, and cost, ten  
53 times higher than cup anemometers, to the long-term site assessment campaigns.

54 For onshore sites there are few studies that analyse how to characterize atmospheric stability and those that exist  
55 are on simple topography in coastal areas (Bardal et al., 2018).

56 Although the behaviour of wind flow over complex terrain is widely studied, as Finnigan summarizes in  
57 (Finnigan et al., 2020) and there are recent publications about the influence of atmospheric stability in wind farms  
58 located in complex terrain (Han et al., 2018; Radünz et al., 2020, 2021); there are no references that analyse in detail  
59 how to characterize atmospheric stability according to different instrumentation requirements.

60 Measuring atmospheric stability in complex terrain has some challenges (compared to flat terrain), one of them  
61 is the fact that the MOST is developed for horizontally homogeneous and flat terrain and in complex terrain vertical  
62 wind speed can be due to stability or sloping terrain, therefore, vertical fluxes will be “contaminated” by terrain  
63 effects. This can be mitigated by using good measurement practices (data quality, coordinate systems and post  
64 processing options) (Stiperski and Rotach, 2015).

65 This study presents atmospheric stability characterization from one mountainous site obtained using two  
66 methods: sonic method and the Richardson bulk number. Measurements of different heights have been used to see  
67 the influence of this parameter on the results

68 The place used in this study meets the characteristics of a typical complex terrain site for wind energy  
69 deployment. The 118 m high MP5 reference meteorological mast, as is explained in other articles by Sanz (Sanz  
70 Rodrigo et al., 2013) and Santos (Santos et al., 2020), is equipped with wind (cup and 3D sonic anemometer) and  
71 temperature measurements distributed along six vertical levels: 2, 40, 80, 90, 100 and 118 m above the ground level  
72 (a.g.l), enabling the comparison between Richardson bulk number and the sonic method to evaluate atmospheric  
73 stability.

74 Special focus is given to explaining the post-processing methodologies to derive stability from raw data  
75 considering fast-response sonic anemometer in a complex terrain.

## 76 2. Atmospheric stability definitions

### 77 2.1 The Obukhov length

78 Monin and Obukhov (M-O) (Monin and Obukhov, 1954) introduced the Obukhov length  $L$  to characterize  
79 atmospheric stability, which is proportional to the height above the surface at which the production of turbulent  
80 energy from buoyancy dominates over mechanical shear production of turbulence (Stull, 1989), and it is defined as:

$$L = -\frac{u_*^3}{\kappa \frac{g}{\Theta_0} \overline{w'\theta'}} \quad (1)$$

81 Where  $g=9.81 \text{ ms}^{-2}$  is the acceleration due gravity,  $\kappa = 0.41$  is the von Karman constant,  $u^*$  is the friction  
82 velocity,  $\Theta_0$  is the surface potential temperature and  $\overline{w'\theta'}$  is the heat flux. The dimensionless height  $\zeta = z/L$  is used  
83 as stability parameter, where  $\zeta < 0$  indicates unstable,  $\zeta > 0$  stable and  $\zeta = 0$  neutral conditions.

84 Table 1 shows the (Sorbjan and Grachev, 2010) stability classification proposing, four regimes in  
85 stable conditions. This classification is also followed by (Sanz Rodrigo et al., 2015) assuming a symmetric  
86 classification in the unstable range. Sanz Rodrigo *et al.* shift the "extremely unstable and stable" regime limit to  $|\zeta| =$   
87 1 in order to avoid contamination of the large scatter found in the high ends of the scale to the "very unstable  
88 and/stable" class. An additional limit is added at  $|\zeta|=0.2$  to give higher resolution in the most frequent stability range.  
89 For consistency, we shall adopt the same classification used in (Sanz Rodrigo et al., 2015) to facilitate the  
90 comparison with offshore conditions.

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**Table 1 Classification of atmospheric stability adapted from (Sorbjan and Grachev, 2010).**

Stability Class	Stability parameter $\zeta = z/L$
extremely unstable (xu)	$\zeta < -1$
very unstable (vu)	$-1 < \zeta < -0.6$
unstable (u)	$-0.6 < \zeta < -0.2$
weakly unstable (wu)	$-0.2 < \zeta < -0.02$
near-neutral (n)	$0 < \zeta < 0.02$
weakly stable (ws)	$0.02 < \zeta < 0.2$
stable (s)	$0.2 < \zeta < 0.6$
very stable (vs)	$0.6 < \zeta < 1$
extremely stable (xs)	$\zeta > 1$

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Using sonic anemometers and eddy covariance technique, the Obukhov length can be obtained. In this way, stability is evaluated locally based on turbulent fluxes averaged over periods from minutes to one hour to integrate the kinetic energy in the microscale turbulence range.

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Sonic anemometer can be used in complex terrain to derive the local Obukhov length. Following the planar fit method of (Wilczak et al., 2001), momentum fluxes should be calculated in the mean streamline plane and heat fluxes in the true vertical coordinate system. If the streamline plane can be known a priori, from a wind direction sector with uniform slope, the planar fit method can be used to infer the mounting tilt angle and correct for it to reduce the uncertainty on the vertical fluxes.

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## 2.2 Bulk Richardson number

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The bulk Richardson number  $Ri_b$  is a form of the Richardson number that is widely used for characterizing stability for its simplicity, defined in terms of a potential temperature difference and a single velocity level:

$$Ri_b = - \frac{gz\Delta\theta}{\theta_0 \overline{U}^2} \quad (2)$$

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Where, as proposed (Sanz Rodrigo et al., 2015), the height  $z$  is taken here as the mean height between the two levels of temperature and  $\Delta\theta$  is derived from the water-air or surface-air temperature difference.

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As (Bardal et al., 2018) propose the general empirical relations from (Businger et al., 1971) slightly modified by (Dyer, 1974) have been used to relate  $\zeta$  with the  $Ri_b$ :

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$$\zeta = \begin{cases} Ri_b, & Ri_b < 0 \\ \frac{Ri_b}{1-5Ri_b}, & 0 < Ri_b < 0.2 \end{cases} \quad (3)$$

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Alternatively  $Ri_b$  can be used directly to do a stability classification. (Mohan, 1998) has proposed a seven classes of stability classification methodology (Table 2) which has been accepted by the scientific community as it was shown by (Ruisi and Bossanyi, 2019).

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**Table 2 Classification of atmospheric stability (Mohan, 1998).**

Stability Class	Stability parameter $Ri_b$
Very unstable	$Ri_b < -0.023$
Unstable	$-0.023 \leq Ri_b < -0.011$
Weakly unstable	$-0.011 \leq Ri_b < -0.0036$
Neutral	$-0.0036 \leq Ri_b < 0.0072$
Weakly stable	$0.0072 \leq Ri_b < 0.042$
Stable	$0.042 \leq Ri_b < 0.084$
Very stable	$Ri_b \geq 0.084$

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## 3. The Alaiz site

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The MP5 mast is located (42°41.7' N, 1°33.5' W) at the top of Alaiz mountain in the region of Navarre (Spain), around 15 km SSE from Pamplona in the CENER's experimental wind farm. The prevailing wind directions are from the North and from the South. To the North there is a large valley at around 700 m lower altitude. To the

121 South, complex terrain is found with the presence of some wind farms; the closest one situated 2 km behind the row  
122 of six wind turbine stands of the test site (see Fig. 1). As it is explained by (Sanz Rodrigo et al., 2013) the wakes  
123 from this wind farm can be considered well mixed with the boundary layer flow in most conditions so additional  
124 turbulence it is not expected in MP5 due wakes from neighbour wind farms.

125 Besides MP5 meteorological mast there are four other reference met masts (MP0, MP1, MP3 and MP6), all of them  
126 118 m tall.

127 The test site started operating in 2009 with the site calibration procedures. The first wind turbines were installed  
128 in the summer of 2011. The standard configuration of each mast is designed for multi-megawatt wind turbine testing  
129 and includes sonic and cup anemometer, wind vanes and temperature/humidity measurements. Replicated cup  
130 anemometers are situated 2 m below the reference ones.

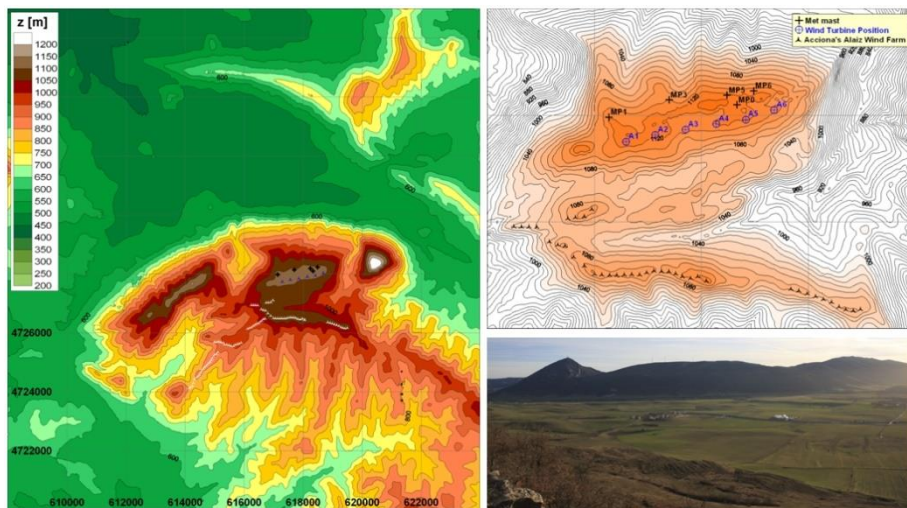
131 The mast MP5 is 118 m high lattice permanent mast with nine measurement levels with booms oriented to the  
132 West (263°) and the East (83°). Wind speed and wind direction are measured at five levels (118, 102, 90, 78 and 40  
133 m) with cups anemometer (oriented to the West) and wind vanes (oriented to the East); while sonic anemometer are  
134 installed at 115.5, 75.5 and 39.5 m (oriented to the West). Temperature and relative humidity are measured at five  
135 levels (113, 97, 81, 38 and 2 m) and pressure at 2 m high.

136 The instrumental set-up is compliant with IEC 61400-12-1(IEC61400-12-1 (ED1) 2005-12, 2005) with  
137 MEASNET cup anemometer calibration (Measnet, 2009) and with ENAC accreditation according to UNE-EN  
138 ISO/IEC 17025.

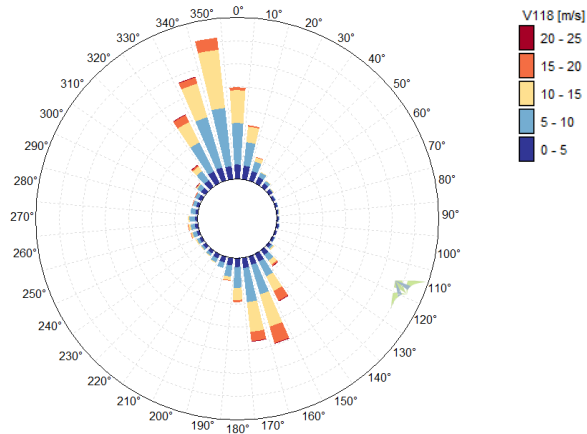
139 The data acquisition system consist in a real-time controller CompactRIO from National Instruments with 128  
140 MB DRAM and 2 GB storage embedded in a chassis in connection with 8 modules of digital and analogical data  
141 acquisition. All connected to an Ethernet network.

142 The rate sample is 5 Hz for cup anemometer (Vector A100LK) and 20 Hz for sonic anemometer (METEK USA-  
143 1), wind vanes (Thies Compact), pressure (Vaisala PTB100A), and humidity temperature sensor (Ammonit P6312).

144 Figure 2 shows the wind rose at the MP5 site, from the period between July 2014 to June 2015. It presents a  
145 bidirectional wind climate, with prevailing winds from the north-northwest sector (330–360, 32% of total) and the  
146 south southeast sector (150–180, 28% of total).



147  
148 **Figure 1 Alaiiz elevation map, close-up of the test site and view from the upstream ridge to the North.**  
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150 **Figure 2 Wind rose of 10 min wind speeds observed by MP5 at 118m for the reference period (July 2014–June 2015).**

151 **4. Methodology**

152 In the present work, a one year period (1st July 2014 to 30th June 2015) is analyzed. Flux measurements from the  
 153 sonic anemometer at 115.5, 75.5 and 39.5 m are used to calculate de Obukhov length  $L$ , while conventional sensors  
 154 (wind direction, relative humidity, air pressure and temperature) are used to estimate the bulk Richardson number.

155 **4.1 Data quality control**

156 Before calculating stability parameter all data are checked for data quality.

157 Data from conventional sensors (wind direction, relative humidity, air pressure and temperature) have been  
 158 processed following Brower (Brower, 2012). It consists on checking the completeness of the collected data and  
 159 applying several test (range, relational and trend). After filtering for quality-control purposes, the conventional  
 160 sensors provide horizontal wind speeds, directions, relative humidity, pressures and temperatures availabilities  
 161 greater than 85% at all levels during the evaluation period.

162 For sonic anemometer there are a lot of procedures (Aubinet et al., 2012) and test criteria for quality control of  
 163 turbulent time series and studies about the impact in the results of this procedures (Stiperski and Rotach, 2015).

164 High-frequency raw data often contain impulse noise, that is, spikes, dropouts, constant values, and noise. Spikes  
 165 in raw data can be caused by instrumental problems, such as imprecise adjustment of the transducers of ultrasonic  
 166 anemometer, insufficient electric power supply, and electronic noise, as well as by water contamination of the  
 167 transducers, bird droppings, cobwebs, etc., or rain drops and snowflakes in the path of the sonic anemometer.

168 Several spikes in wind speed have been detected in the raw sonic anemometer data. Therefore, a de-spiking filter  
 169 is applied based on the change in wind speed from each data point to the next and taking into account the physical  
 170 limits according to sensor specifications. Data points are removed if they are preceded and followed by changes  
 171 exceeding the lowest 99% of all changes. After filtering the spikes, the sonic anemometer provide wind speed and  
 172 temperature availabilities greater than 80% in the three sonic anemometer.

173 **4.2 Eddy Covariance method**

174 The operating principles of sonic anemometer are described by different authors (Aubinet et al., 2012; Cuerva et al.,  
 175 2003; Kaimal and Businger, 1963; Kaimal and Finnigan, 1994; Schotanus et al., 1983). The sonic anemometer  
 176 output provides three wind components in an orthogonal axis system and sonic temperature. The relation between  
 177 sonic temperature and absolute real temperature is given by (Kaimal and Gaynor, 1991).

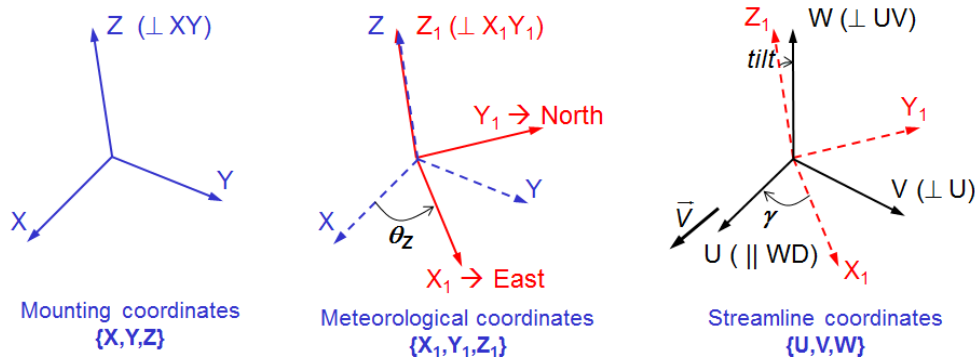
178 High frequency data from sonic anemometer have been processing to obtain 10 minutes databases that include  
 179 turbulent fluxes of energy, mass, and momentum with the eddy covariance technique (Aubinet et al., 2012)(Burba,  
 180 2013; Burba and Anderson, 2010; Geissbühler et al., 2000).

181 The main requirements for instruments and data acquisition systems used for eddy covariance data are their  
 182 response time to solve fluctuations up to 10 Hz. This means that the sampling frequency has to be high enough to

183 cover the full range of frequencies carrying the turbulent flux, leading usually to a sampling rate of 10–20 Hz. In the  
 184 test case in this report 20 Hz is the sample rate for the sonic anemometer.

185 The transformation of high-frequency signals into means, variances, and covariances requires different steps  
 186 (Aubinet et al., 2012; Stiperski and Rotach, 2015), in this study the next steps have been proposed:

- 187 1. Quality Control of raw data, explained in point 4.1.2
- 188 2. Coordinate Rotation, transformation of coordinate systems, from the original axes based on the anemometer  
 189 output to the streamline terrain-following system, based on the Planar Fit Method (PFT) (Richiandone et al.,  
 190 2008; Wilczak et al., 2001). Momentum fluxes and heat fluxes have been calculated with respect to the  
 191 streamline terrain-following coordinate system. Figure 3 shows the steps to rotate the axes from mounting  
 192 coordinates to streamline coordinates.



193 **Figure 3 Schematic description for the rotation process.**

- 194 3. Variance and Covariance Computation, apply eddy covariance technique for calculation of vertical turbulent  
 195 fluxes (heat and momentum). It corresponds to the calculation of the covariance of the fluctuations of the vertical  
 196 velocity with the quantity  $\Phi$  (temperature for heat, velocity components for momentum).

$$F_{\phi} = \overline{w'\phi'} = \overline{w\phi} - \overline{w}\overline{\phi} = \frac{1}{N-1} \left[ \sum w'\phi' - \frac{1}{N} (\sum w') (\sum \phi') \right] \quad (4)$$

197 N denotes the number of samples considered for the short averaging period T over which the flux is  
 198 calculated (from 5 to 60 minutes). N has to be long enough to ensure statistical convergence and short enough to  
 199 assume stationarity (in complex terrain difficult to fulfil both criteria). In this work a 10 minutes averaging period  
 200 has been selected.  
 201

202 In the MP5's sonic anemometer, at 115.5, 75.5 and 39.5 m height, moreover the temperatures, the variables  
 203 recorded are: the module of wind speed vector, the direction and vertical component (z). These values are projected  
 204 to meteorological coordinates to obtain the three components of wind speed vector (x, y, z) after being filtered the  
 205 transformation of high-frequency signals into means, variances, and covariances has been done.

206 The 10 minutes values of wind speed from sonic anemometer after applying steps 1 to 3 are checked and some  
 207 non-valid data are detected. As in conventional sensors these invalid data are due to icing effects so they are filtered.

## 208 4.4 Stability assessment

209 MP5's sonic anemometer allow evaluating stability based on the local Obukhov length at different heights. This will  
 210 be the benchmark method since it is directly obtained from the measurements without introducing any assumptions  
 211 or empirical relationships. The bulk Richardson number is evaluated as an alternative methodology since it follows  
 212 easier instrumentation set-up and post-processing, and for offshore places has presented good results (Sanz Rodrigo,  
 213 2011; Sanz Rodrigo et al., 2015).

### 214 4.4.1 Sonic method

215 To obtain the stability parameter  $\zeta = z/L$ , as it was explained before, sonic anemometer measurements are rotated to  
 216 the mean streamline coordinate system using the planar fit method to guarantee that the mean streamline plane will  
 217 be parallel to the terrain surface. After this, variances and covariances of detrended velocity and sonic temperature  
 218 perturbations are computed using the eddy covariance technique over high frequency timescale. Then, turbulent  
 219 fluxes are obtained by averaging the covariances over a period of 10 minutes.

220 In complex terrain, the hypothesis of a homogeneously horizontal surface layer is not fulfilled so the applicability of  
 221 Monin and Obukhov similarity theory (MOST) to complex terrain conditions is not obvious. This signify that for the  
 222 complex sites as Alaiz the theory is not completely valid because the topography creates local variations of wind  
 223 flow near the ground (Kaimal and Finnigan, 1994).

#### 224 4.4.2 Bulk Richardson number

225 As it was explained before, sonic anemometry is not routinely used in wind energy, and bulk Richardson number  $Ri_b$   
 226 is an alternative way to estimate atmospheric stability based on a temperature difference and a single velocity level  
 227 (Eq. (2)).

228 In  $Ri_b$  number equation, potential temperature  $\Theta$ , is the temperature of an air parcel with absolute temperature  $T$   
 229 and pressure  $p$  would have if brought adiabatically to the pressure at the 1000 mb level. To first order it can be  
 230 calculated as:

$$\theta = T + \left(\frac{g}{C_p}\right) \Delta z \quad (5)$$

231 Where  $g$  is the acceleration due gravity,  $C_p=1000 \text{ Jkg}^{-1}\text{K}^{-1}$  is the specific heat capacity of the air at constant  
 232 pressure, and  $\Delta z$  is the height difference from the 1000 mb level.

233 With Eq. (3) the obtained  $Ri_b$  will be used to estimates the stability parameter  $\zeta = z/L$ . As Bardal *et al.* (Bardal *et*  
 234 *al.*, 2018) explain, these formulations are only valid for values lower than 0.2, but to make a classification according  
 235 to atmospheric stability they are considered adequate.

### 236 5. Results and discussion

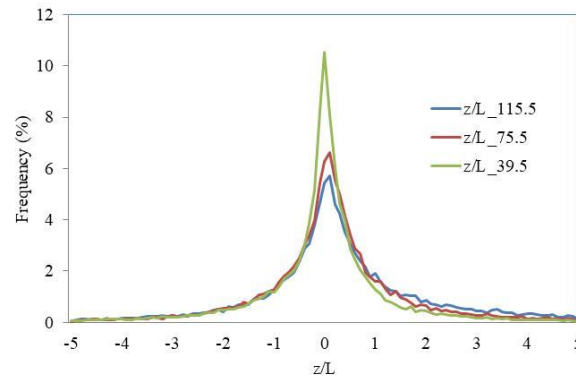
237 The study is divided into two parts: statistics of atmospheric stability with both methods (the Obukhov length and  
 238 Richardson Bulk); and comparison between both methods.

#### 239 5.1 Sonic method

240 Atmospheric boundary layer (ABL) models used in wind farm design tools are typically based on Monin-Obukhov  
 241 theory. In stable conditions this surface-layer theory is extended to the entire ABL by assuming local scaling of  
 242 turbulence characteristics through the stability parameter  $\zeta = z/L$ . This similarity theory would produce self-similar  
 243 profiles of dimensionless quantities regardless of the height above ground level.

244 In the study case, as it was explained before, from the high-frequency (20 Hz) data recorded in the three  
 245 available sonic anemometers in MP5 mast, the values of the Obukhov length ( $L$ ) over a period of 10 minutes have  
 246 been obtained, and taking into account the heights at which they are installed, the parameter  $\zeta = z/L$ .

247 In Fig. 4 the stability parameter  $\zeta = z/L$  frequency distribution at the three sonic heights is depicted, resulting in  
 248 showing a good agreement among them with a reduction of the percentage of conditions near neutral stability as the  
 249 measurement height increases.



250 **Figure 4 Probability distribution of  $z/L$  at all the sonic heights. Only concurrent time steps between July 2014 and June**  
 251 **2015 are included.**  
 252

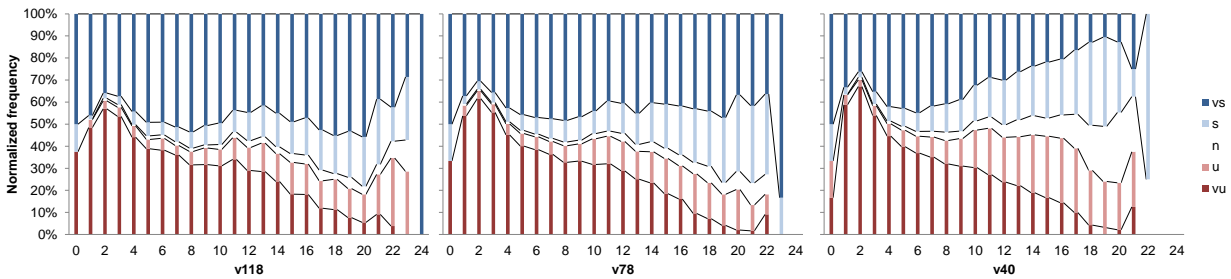
253 Figure 5 shows the distribution of atmospheric stability against wind speed at the MP5 measurements heights,  
 254 the 9 stability classes propose in Table 1 are reduced to five combining: weakly unstable and stable classes with

255 unstable and stable classes; and very unstable and stable with extremely unstable and stable. Table 3 shows the  
 256 classification used. For the three heights, the stable situations are slightly higher than the unstable ones and there is  
 257 an increase of neutral and stable conditions with increasing wind speeds, this is in accordance with the general  
 258 knowledge that for strong wind speeds the atmosphere becomes neutrally stratified.

259 **Table 3 According to Table 1 a reduced five stability classes.**

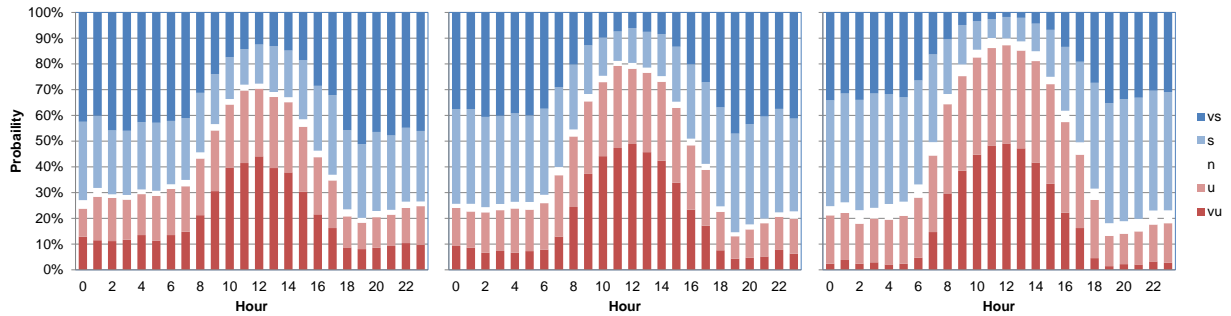
Stability Class	Stability parameter $\zeta = z/L$
very unstable (vu)	$-\zeta < -0.6$
unstable (wu)	$-0.6 < \zeta < -0.02$
neutral (n)	$-0.02 < \zeta < 0.02$
stable (ws)	$0.02 < \zeta < 0.6$
very stable (vs)	$0.6 < \zeta$

260  
 261 As mentioned before, it is observed a significant dependence of stability distributions with height. At higher  
 262 levels, the stability distributions are broader and there are more frequent cases with very large and extreme stability.  
 263 This dependency of the stability distribution with height is because  $z$  is part of the definition of the stability  
 264 parameter; and closer to the ground there are more “neutral” conditions because  $z/L$  tends to zero.



265 **Figure 5 Distribution of atmospheric stability with wind speed based on  $z/L$  obtained with sonic anemometer at different**  
 266 **heights, 115.5 m on the left, 75.5 m in the middle and 39.5 m on the right side. vs, very stable; s, stable; n, neutral; u,**  
 267 **unstable; vu, very unstable.**

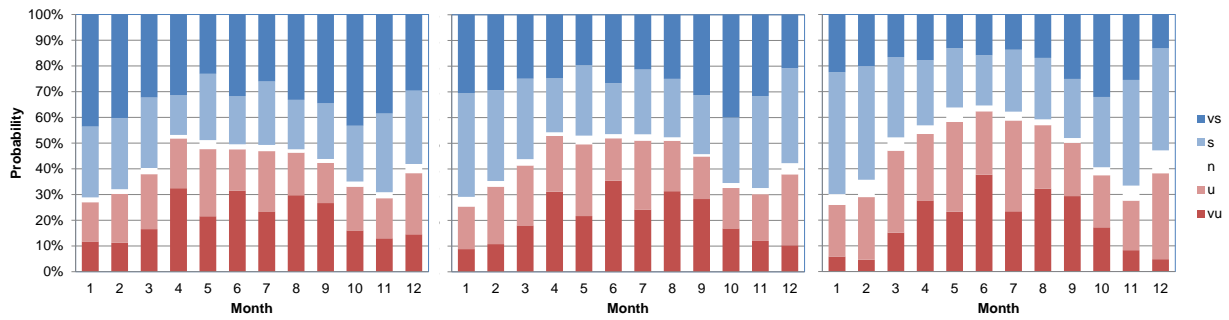
268 The diurnal cycle, see Fig. 6, presents unstable conditions developing from 9.00 to 15.00. The rest of the day is  
 269 dominated by stable conditions resulting in low turbulence intensities.



270 **Figure 6 Distribution of atmospheric stability with hour based on  $z/L$  obtained with sonic anemometer at different**  
 271 **heights, left 115.5 m, center 75.5m and right 39.5 m. vs, very stable; s, stable; n, neutral; u, unstable; vu, very unstable.**

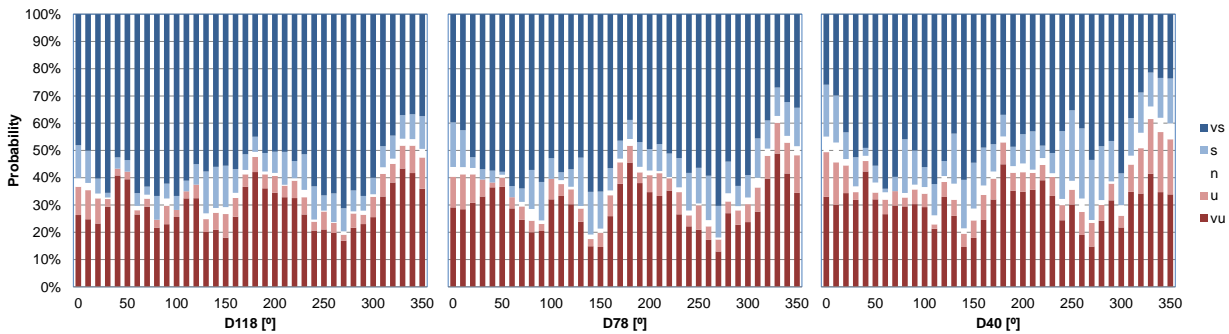
272  
 273 Figure 7 shows the evolution of stability throughout the year. The stable side dominates during winter months,  
 274 with unstable conditions peaking between April to August where they take a  $\approx 50\%$  share.





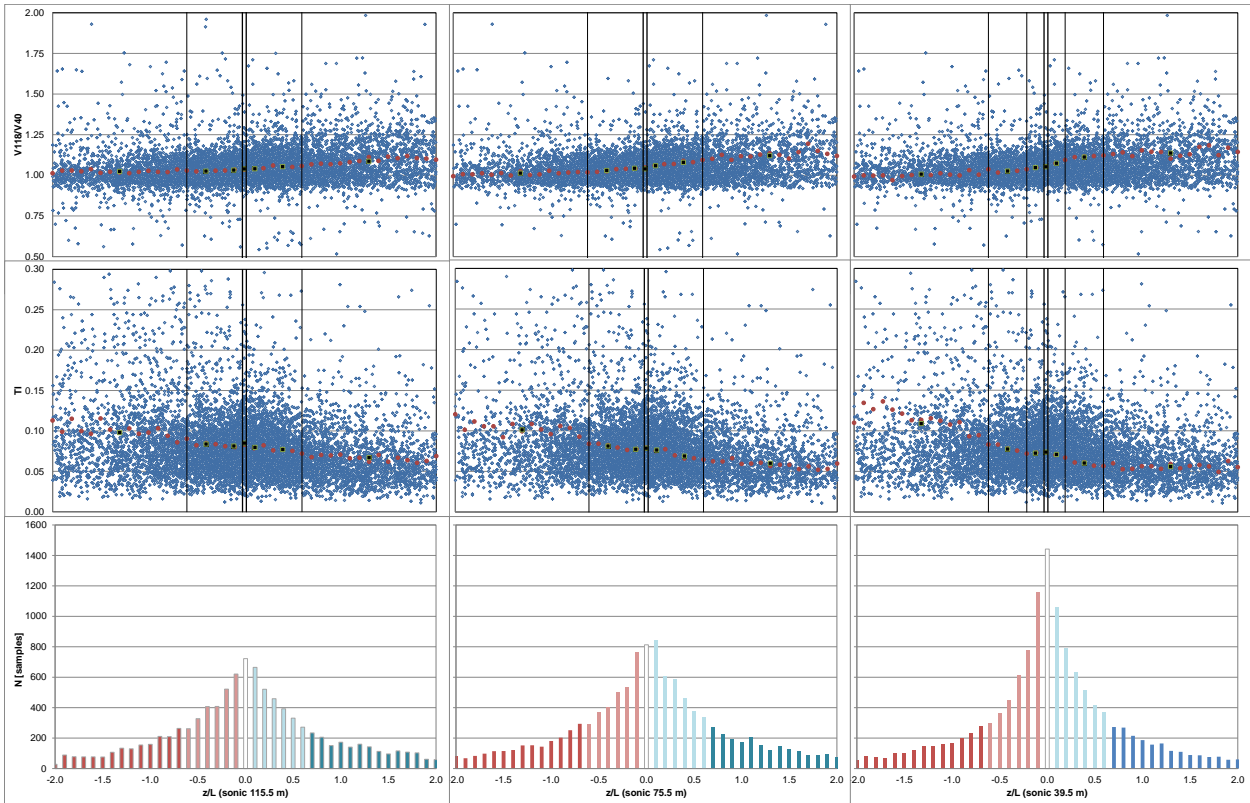
275 **Figure 7 Monthly distribution of stability based on  $z/L$  obtained with sonic anemometer at different heights, left 115.5**  
 276 **m, center 75.5 m and right 39.5 m. vs, very stable; s, stable; n, neutral; u, unstable; vu, very unstable.**

277 The variation of atmospheric stability with wind direction is showed in Fig. 8. Stable situations dominate in most  
 278 of the directions except for the northwest direction ( $330^{\circ}$ - $350^{\circ}$ ) that is one of the predominant in Alaiz. As can be  
 279 seen in Fig. 1, the North face of Alaiz Mountain has a steep slope (the Roughness Index (RIX) value in the north  
 280 sector in MP5 position is 22.4%) that empties into a large valley at around 700 m lower altitude. According to (Stull,  
 281 1989) this topography causes ascending hillside/valley winds that generate convective turbulence and therefore  
 282 situations of instability that could explain some of the unstable conditions found in the  $330^{\circ}$ - $350^{\circ}$  direction sector.

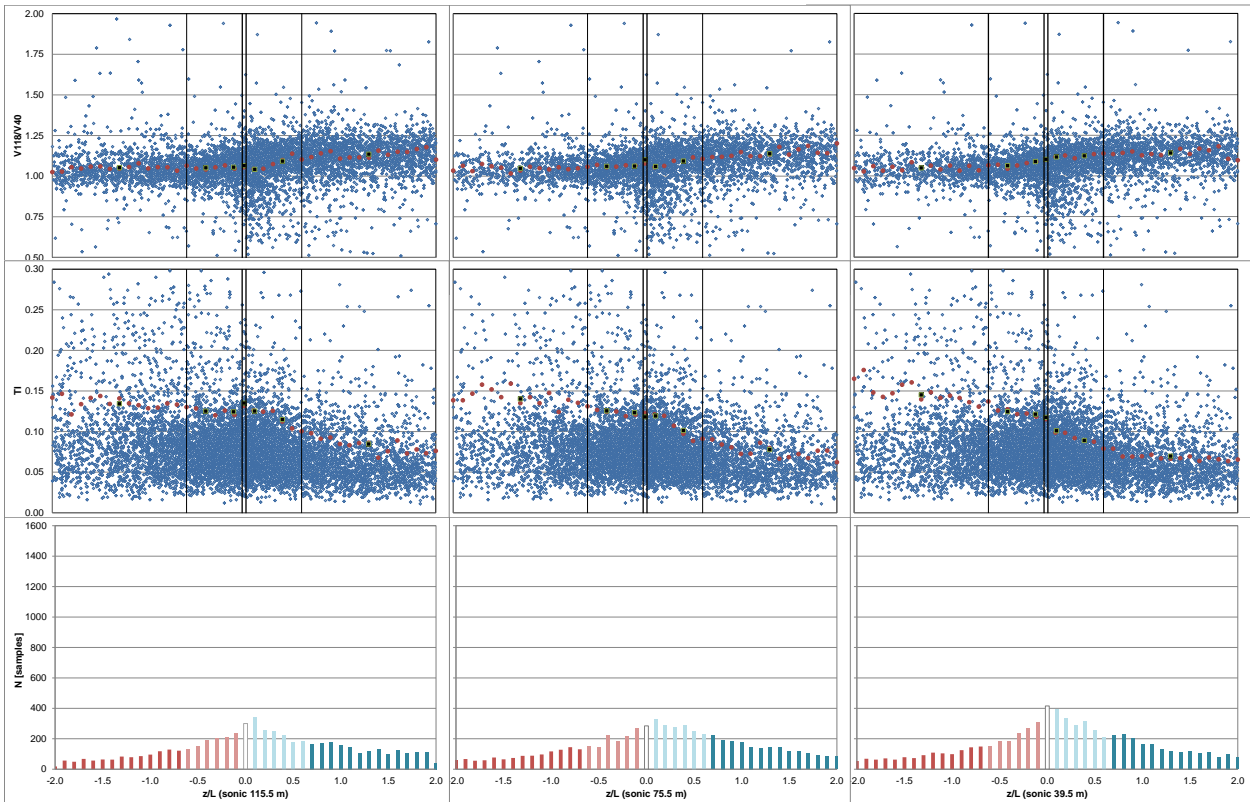


283 **Figure 8 Distribution of atmospheric stability with wind direction based on  $z/L$  obtained with sonic anemometer at**  
 284 **different heights. vs, very stable; s, stable; n, neutral; u, unstable; vu, very unstable.**

285 Following the stability classification defined in Table 3, Fig. 9 and Fig. 10 present the dependency of wind shear  
 286 (calculated as the wind speed ratio between 118 and 40 m) and turbulence intensity (calculated as the ratio of the  
 287 standard deviation to the mean wind speed at 118 m) with stability based on  $z/L$  parameter from the three sonic  
 288 sensors installed for the NNW and SSE prevailing wind direction sectors.  
 289



290 **Figure 9** Wind shear and turbulence intensity vs sonic stability in *MP5*, [337.5°-22.5°] sector. Red dots are the  $z/L$  mean  
 291 values for 0.01 resolution scale, black squares are the  $z/L$  mean values in each of the stability classes according to Table 1.  
 292



293 **Figure 10** Wind shear and turbulence intensity vs sonic stability in *MP5*, [157.5°-202.5°] sector. Red dots are the  $z/L$  mean  
 294 values for 0.01 resolution scale, black squares are the  $z/L$  mean values in each of the stability classes according to Table 1.

295 For the three heights is observed that, as is explained by (Emeis, 2013), in unstable situations the ground surface is  
296 warmer than the air above so there is a positive heat flux that causes more turbulence. This results in a convective,  
297 well-mixed, surface layer with small vertical gradients. On the other hand, lower turbulence and high shear wind  
298 profiles are associated to stable situations where turbulence is reduced due to a negative vertical heat flux.

## 299 5.2 Bulk Richardson number

300 Since sonic anemometers are not commonly used in wind resource assessment, an alternative method to estimate the  
301 atmospheric stability is Bulk Richardson number. It is based on mean wind speed at height  $z$  and mean virtual  
302 potential temperature difference between air at the reference height ( $z$ ) and surface temperature.

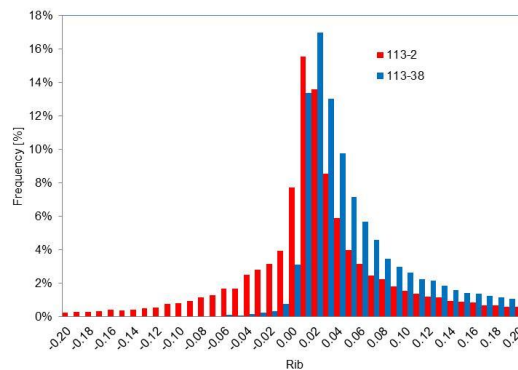
303 The calculation of the Bulk–Richardson number is, in the present study, not straightforward because of the lack  
304 of reliable sensors at the surface. The lower air temperature is measured at 2 m in MP5 mast. Ideally, the  
305 temperature difference at the air-surface interface is required (Kaimal and Finnigan, 1994) for stability analysis.  
306 However, because of the lack of surface temperature, 2 m height air temperature has been chosen as representative.  
307 Observations of 118 m wind speed and 113 m air temperature have been used in conjunction with 2 m air  
308 temperature to estimate  $Ri_b$ .

309 The MP5 mast has not measurements of surface temperature or near the ground. Some authors in these  
310 circumstances either extrapolate the values to the surface ( $z=0$ ) (Machefaux et al., 2016) or perform the calculation  
311 directly between the available temperature levels (Martin et al., 2016; Ruisi and Bossanyi, 2019; Zhan et al., 2020).  
312 To analyze how the choice of measurement heights may influence resulting  $Ri_b$  stability distributions the  $Ri_b$  has  
313 also been calculated using 38 m air temperature instead 2 m.

314 The values of the Bulk-Richardson number have been obtained over a period of 10 minutes, i.e. the same period  
315 used for calculation of the Obukhov length.

316 Figure 11 shows the distribution for the bulk Richardson number method. The lower measurement level is varied  
317 between 2 and 38 m. Using the 38 m level, it is observed that according to the classification in Table 2, unstable  
318 cases practically disappear. This is not physically possible and does not occur in the classification obtained by the  
319 sonic method (see Fig. 4). So In this case, the results obtained using the 38 m temperature sensor as a representative  
320 surface level do not give us any reliable information. Small temperature differences highly affect the result of the  
321 Richardson number method and therefore it is greatly affected by deviations in the measurement of this variable.  
322 The MP5 temperature sensors have an accuracy of 0.3 °C and the mean temperature difference in the period  
323 analyzed between the level of 38 m and that of 113 m has been 0.7 °C so the uncertainty of the measurement is of  
324 the same order as the measurement itself.

325 The selection of temperature measurement heights has a great effect on the bulk Richardson number method,  
326 both in the exactitude and in the applicability of the method. To reduce uncertainties the measurements should be  
327 made either with differential temperature sensors or with calibrated sensors and a sufficient vertical separation in  
328 order to reduce the influence of inaccuracies in the temperature measurements (Baker and Bowen, 1989; Brower,  
329 2012).



330  
331 **Figure 11 Probability distribution of  $Ri_b$  measured between 2 m and 113 (red one) and between 38 and 113 m (blue lines).**  
332 **Only concurrent time steps between July 2014 and June 2015 are included.**

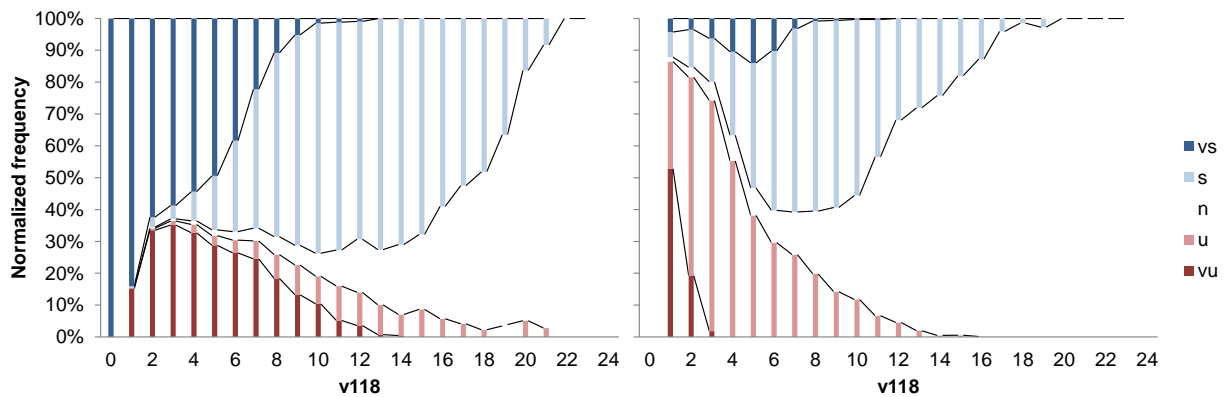
333  
334 Figure 12 shows the distribution of atmospheric stability against wind speed. On the left side atmospheric  
335 stability is directly classified with the  $Ri_b$  obtained with observations of 118 m wind speed, 113 m air temperature  
336 and 2 m air temperature, this last temperature sensor has been chosen as representative of surface temperature. The  
337 seven stability classes propose in Table 2 are reduced to five combining: weakly unstable and stable classes with

338 unstable and stable classes. Table 4 shows the classification used. On the right side atmospheric stability is classified  
 339 according to the stability parameter  $\zeta = z/L$  obtained with  $Ri_b$  and Eq. (3) and according to stability classification  
 340 proposed in Table 3.

341 **Table 4 According to Table 2 a reduced five stability classes.**

Stability Class	Stability parameter $Ri_b$
Very unstable	$Ri_b < -0.023$
Unstable	$-0.023 \leq Ri_b < -0.0036$
Neutral	$-0.0036 \leq Ri_b < 0.0072$
Stable	$0.0072 \leq Ri_b < 0.084$
Very stable	$Ri_b \geq 0.084$

343  
 344 Both distributions show a differentiated behavior with fewer “very” unstable and stable situations and a greater  
 345 number of neutral observations in the case of the classification with  $\zeta$  (on the right side of Fig. 12).



346 **Figure 12 Distribution of atmospheric stability with wind speed. On the left based on  $Ri_b$ ; On the right based on  $z/L$**   
 347 **obtained from  $Ri_b$  with transformation functions by Businger and Dyer. vs, very stable; s, stable; n, neutral; u,**  
 348 **unstable; vu, very unstable.**

### 349 5.3 Comparison of stability methods: sonic versus bulk method

350 Comparing the distribution of atmospheric stability against wind speed based on sonic method (Fig. 5) with the  
 351 results obtained based on  $Ri_b$  method (Fig. 12); it is observed that there are important differences between them.

352 Table 5 presents a frequency of occurrence of stability classes with concurrent data using different methods. This  
 353 quantitative comparison shows that taking the sonic method as benchmark, it is observed that the bulk method when  
 354 the Businger and Dyer functions are used to estimate the stability parameter  $\zeta = z/L$  over predict the percentage of  
 355 neutral and stable conditions to the detriment of very unstable and stable situations, probably due to similar air  
 356 temperature values at 113 and 2 m. On the other hand, classification directly with  $Ri_b$  according to Mohan  
 357 classification over predict too the stable situations at the cost of under predicting the unstable ones. As is explained  
 358 in some references (Bardal et al., 2018; Sathe et al., 2011), stability characterization with  $Ri_b$  have several weak  
 359 points: in one hand  $Ri_b$  method is sensitive to temperature measurements and uncertainty in  $L$  estimation increases  
 360 as the temperature difference is reduced. Besides, other source of uncertainty comes from the definition of the  
 361 surface temperature. In the other hand Businger and Dyer functions have some limitations and as (Bardal et al.,  
 362 2018) propose the use of more advanced methods for relating the  $Ri_b$  to the  $z/L$  parameter might improve the results.

363 Besides these methodological reasons there are some physical causes of the differences found. One of these is  
 364 that Richardson bulk number represents a bulk average stability value instead a local measurement like the sonic  
 365 method.

366 **Table 5 Frequency of occurrence of stability classes.**

	115.5/L	75.5/L	39.5/L	$z/L$ from $Ri_b$	$Ri_b$
vu	21.2%	21.3%	19.9%	0.7%	18.1%
u	19.4%	21.4%	26.8%	21.2%	5.9%
n	2.2%	2.4%	4.4%	32.5%	8.2%

s	24.0%	28.2%	29.9%	42.2%	43.6%
vs	33.2%	26.7%	19.1%	3.5%	24.2%

## 367 6. Conclusions

368 In this work, a detailed data analysis focused on how to estimate atmospheric stability in a site with complex terrain  
369 was presented. The Obukhov parameter  $\zeta = z/L$ , which can be measured locally with the use of a sonic anemometer,  
370 and bulk Richardson number have been studied. The methods are examined considering their theoretical  
371 background, implementation complexity, instrumentation requirements, and practical use in connection with wind  
372 energy applications.

373 It is shown that the resulting stability depends on which method is chosen. The sonic method is taken as  
374 benchmark because is the only way of measuring local stability without the use of empirical functions or theoretical  
375 assumptions. However this method requires working with accurate high frequency data, rotating the measurements  
376 to align the coordinate system to the mean wind vector, which is reported to require special attention in complex  
377 terrain to guarantee that the mean streamline plane will be parallel to the terrain surface; to finally obtain turbulent  
378 fluxes using the eddy covariance technique.

379 According to the stability parameter  $\zeta = z/L$  obtained with the three sonic anemometer installed in MP5 mast.  
380 For the three heights, the stable situations are slightly higher than the unstable ones and there is an increase of  
381 neutral and stable conditions with increasing wind speeds. There is a significant dependence of stability distributions  
382 with height. At higher levels, the stability distributions are broader and there are more frequent cases with very large  
383 and extreme stability.

384 The seasonal and diurnal cycle is identified, in the winter and during the hours between 17h to 8h stable side  
385 dominates, while between April to August and between 9h to 15h unstable conditions are found to be more frequent.  
386 Winds from the predominant northwest direction (330°-350°) produce more unstable conditions than the others  
387 sectors.

388 For the three heights, and in the two predominant sectors, is observed that in in unstable situations the ground  
389 surface is warmer than the air above so there is a positive heat flux that causes more turbulence. This results in a  
390 convective, well-mixed, surface layer with small vertical gradients. On the other hand, lower turbulence and high  
391 shear wind profiles are associated to stable situations where turbulence is reduced due to a negative vertical heat  
392 flux.

393 As alternative to characterize stability, the bulk Richardson number is explored, it requires the minimum level of  
394 instrumentation, mean wind speed at height  $z$  and mean virtual potential temperature difference between air at the  
395 reference height ( $z$ ) and surface temperature. The bulk Richardson number can be used directly to classify the  
396 atmospheric stability or it can be transform into  $\zeta = z/L$  by Businger and Dyer functions.

397 On MP5 there is not a surface temperature sensor so 2 m high air temperature sensor has been chosen as  
398 representative, moreover to analyze how the choice of measurement heights may influence resulting  $Ri_b$  stability  
399 distributions, it has also been calculated using 38 m air temperature sensor instead 2 m. This configuration does not  
400 give us any reliable information, it could be due temperature sensors on MP5 have an accuracy of 0.3°C and the  
401 mean temperature difference in the period analyzed between the level of 38 m and that of 113 m has been 0.7 °C so  
402 the uncertainty of the measurement is of the same order as the measurement itself. The  $Ri_b$  number relies on smaller  
403 temperature differences for estimation of the mean gradient and its accuracy is therefore dependent on the sensor  
404 precision, calibration and measurement heights. The following recommendations are provided to obtain consistent  
405 results with bulk Richardson method: use high precision temperature sensors; calibrate all the temperature sensors at  
406 the same time; calibrate the temperature sensors in the operational range to guarantee better calibration in the  
407 temperatures of interest and have a reference temperature sensor below 2 m, as close to the ground as possible. On  
408 the other hand, the stability classification obtained using directly the  $Ri_b$  values shows a differentiated behavior than  
409 that estimated according to the stability parameter  $\zeta = z/L$  obtained with  $Ri_b$  and Businger and Dyer functions. It  
410 could be by the different classification employed in both characterization (Mohan vs Sorbjan & Grachev) and/or by  
411 the Businger and Dyer functions.

412 In summary the sonic method is more costly and complex but, in this study, it shows results in accordance with  
413 the general atmospheric boundary layer knowledge, so we recommend it at as first option to obtain a local  
414 measurement of atmospheric stability that can be associated to a certain height above the ground. For the Bulk  
415 Richardson number, based in the references read, there is no a standard methodology for characterizing atmospheric  
416 stability using this method and there are many different approximations. Furthermore, empirical relations to relate  
417  $Ri_b$  to  $\zeta = z/L$  are obtained either for offshore sites or for non-complex sites, so there is a need for observational

418 studies on complex terrain to increase under-standing of how estimate atmospheric stability accurately with Bulk  
419 Richardson number.

420 **Data availability.** Data belongs to CENER and it can be obtained from the author upon request.

421 **Author contribution.** EC is the principal investigator of the project and coordinated the activities and the  
422 preparation of the paper. DP aided in the formulation of the scope of the work, FB assisted in the measurement post-  
423 processing, while the methodology was devised by EC, JS and DP. The stability analysis and visualization was  
424 performed by EC. EC wrote the original draft, AG helped with the composition of the manuscript while EC, JSR,  
425 FB, DP and AG contributed, reviewed and edited the final paper.

426 **Competing interests.** The authors declare that they have no conflict of interest.

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## 429 References

430 Abkar, M. and Porté-Agel, F.: Influence of atmospheric stability on wind-turbine wakes: A large-eddy simulation  
431 study, *Phys. Fluids*, 27(3), 35104, doi:10.1063/1.4913695, 2015.

432 Aubinet, M., Vesala, T. and Papale, D.: *Eddy Covariance : A practical guide to measurement and data analysis.*,  
433 2012.

434 Baker, B. and Bowen, J.: *Quality Assurance Handbook for Air Pollution Measurement Systems: Volume IV :*  
435 *Meteorological Measurements.*, 1989.

436 Bardal, L. M., Onstad, A. E., Sætran, L. R. and Lund, J. A.: Evaluation of methods for estimating atmospheric  
437 stability at two coastal sites, *Wind Eng.*, 42(6), 561–575, doi:10.1177/0309524X18780378, 2018.

438 Brower, M. C.: *WIND RESOURCE ASSESSMENT: A Practical Guide to Developing a Wind Project.*, 2012.

439 Burba, G.: *Eddy Covariance Method for Scientific, Industrial, Agricultural and Regulatory Applications.*, 2013.

440 Burba, G. and Anderson, D.: *Eddy Covariance Flux Measurements.* [online] Available from:  
441 <http://www.ncbi.nlm.nih.gov/pubmed/18767616>, 2010.

442 Cuerva, A., Sanz-Andrés, A. and Navarro, J.: On multiple-path sonic anemometer measurement theory, *Exp. Fluids*,  
443 34(3), 345–357, doi:10.1007/s00348-002-0565-x, 2003.

444 Cuerva, A., Sanz-Andrés, A., Franchini, S., Eecen, P., Busche, P., Pedersen, T. F. and Mouzakis, F.: *ACCUWIND:*  
445 *Task 2. Improve the Accuracy of Sonic Anemometers.*, 2006.

446 Dyer, A. J.: A review of flux-profile relationships, *Boundary-Layer Meteorol.*, 7(3), 363–372,  
447 doi:10.1007/BF00240838, 1974.

448 Emeis, S.: *Wind Energy Meteorology Atmospheric Physics for Wind Power Generation (Green Energy and*  
449 *Technology)*, First Edit., Springer, Berlin, Heidelberg., 2013.

450 Finnigan, J., Ayotte, K., Harman, I., Katul, G., Oldroyd, H., Patton, E., Poggi, D., Ross, A. and Taylor, P.:  
451 *Boundary-Layer Flow Over Complex Topography*, *Boundary-Layer Meteorol.*, 177(2–3), 247–313,  
452 doi:10.1007/s10546-020-00564-3, 2020.

453 Foken, T.: 50 years of the Monin-Obukhov similarity theory, *Boundary-Layer Meteorol.*, 119(3), 431–447,  
454 doi:10.1007/s10546-006-9048-6, 2006.

455 Geissbühler, P., Siegwolf, R. and Eugster, W.: *Eddy Covariance Measurements On Mountain Slopes: The*  
456 *Advantage Of Surface-Normal Sensor Orientation Over A Vertical Set-Up*, *Boundary-Layer Meteorol.*, 96(3), 371–  
457 392, doi:10.1023/A:1002660521017, 2000.

458 Grachev, A. A. and Fairall, C. W.: Dependence of the M-O stability parameters on the Bulk Ri over the ocean, *J.*  
459 *Appl. Meteorol.*, 36, 406–415, doi:10.1175/1520-0450(1997)036<0406:DOTMOS>2.0.CO;2, 1997.

460 Han, X., Liu, D., Xu, C. and Shen, W. Z.: Atmospheric stability and topography effects on wind turbine  
461 performance and wake properties in complex terrain, *Renew. Energy*, 126, 640–651,  
462 doi:10.1016/j.renene.2018.03.048, 2018.

463 Hansen, K. H., Barthelmie, R. J., Jensen, L. E. and Sommer, A.: The impact of turbulence intensity and atmospheric  
464 stability on power deficits due to wind turbine wakes at Horns Rev wind farm., *Wind Energy*, 15, 183–196, 2010.

465 IEC61400-1 (ED4) 2019: IEC61400-1 ed.4, Wind energy generation systems - Part 1: Design requirements. [online]  
466 Available from: <https://webstore.iec.ch/publication/26423> (Accessed 16 April 2021), 2019.

467 IEC61400-12-1 (ED1) 2005-12: IEC61400-12-1 ed1., Wind turbines- Part 12-1: Power performance measurements  
468 of electricity producing wind turbines., 2005.

469 Kaimal, J. C. and Businger, J. a.: A Continuous Wave Sonic Anemometer-Thermometer, *J. Appl. Meteorol.*, 2(1),  
470 156–164, doi:10.1175/1520-0450(1963)002<0156:ACWSAT>2.0.CO;2, 1963.

471 Kaimal, J. C. and Finnigan, J. J.: *Atmospheric Boundary Layer Flows: Their Structure and Measurement*, 1st ed.,  
472 Oxford University Press., 1994.

473 Kaimal, J. C. and Gaynor, J. E.: Another look at sonic thermometry, *Boundary-Layer Meteorol.*, 56(4), 401–410,  
474 doi:10.1007/BF00119215, 1991.

475 Kelly, M., Larsen, G., Dimitrov, N. K. and Natarajan, A.: Probabilistic Meteorological Characterization for Turbine  
476 Loads, *J. Phys. Conf. Ser.*, 524, 012076, doi:10.1088/1742-6596/524/1/012076, 2014.

477 Lange, B., Larsen, S., Højstrup, J. and Barthelmie, R.: Importance of thermal effects and sea surface roughness for  
478 offshore wind resource assessment, *J. Wind Eng. Ind. Aerodyn.*, 92(11), 959–988, doi:10.1016/j.jweia.2004.05.005,  
479 2004a.

480 Lange, B., Larsen, S., Højstrup, J. and Barthelmie, R.: The Influence of Thermal Effects on the Wind Speed Profile  
481 of the Coastal Marine Boundary Layer, *Boundary-Layer Meteorol.*, 112(3), 587–617,  
482 doi:10.1023/B:BOUN.0000030652.20894.83, 2004b.

483 Machefaux, E., Larsen, G. C., Koblitz, T., Troldborg, N., Kelly, M. C., Chougule, A., Hansen, K. S. and Rodrigo, J.  
484 S.: An experimental and numerical study of the atmospheric stability impact on wind turbine wakes, *Wind Energy*,  
485 19(10), 1785–1805, doi:<https://doi.org/10.1002/we.1950>, 2016.

486 Martin, C. M. S., Lundquist, J. K., Clifton, A., Poulos, G. S. and Schreck, S. J.: Wind turbine power production and  
487 annual energy production depend on atmospheric stability and turbulence, *Wind Energy Sci.*, 1(2), 221–236,  
488 doi:10.5194/wes-1-221-2016, 2016.

489 Measnet: MEASNET Procedure: Anemometer Calibration, Version 2, October 2009. [online] Available from:  
490 [http://www.measnet.com/wp-content/uploads/2011/06/measnet\\_anemometer\\_calibration\\_v2\\_oct\\_2009.pdf](http://www.measnet.com/wp-content/uploads/2011/06/measnet_anemometer_calibration_v2_oct_2009.pdf), 2009.

491 Mohan, M.: Analysis of various schemes for the estimation of atmospheric stability classification, *Atmos. Environ.*,  
492 32(21), 3775–3781, doi:10.1016/S1352-2310(98)00109-5, 1998.

493 Monin, A. S. and Obukhov, A. M.: Basic laws of turbulent mixing in the atmosphere near the ground, *Tr. Akad.*  
494 *Nauk SSSR Geofiz. Inst.*, 24, 163–187, 1954.

495 Peña, A. and Hahmann, A. N.: Atmospheric stability and turbulence fluxes at Horns Rev—an intercomparison of  
496 sonic, bulk and WRF model data, *Wind Energy*, 15(5), 717–731, doi:<https://doi.org/10.1002/we.500>, 2012.

497 Radünz, W. C., Sakagami, Y., Haas, R., Petry, A. P., Passos, J. C., Miqueletti, M. and Dias, E.: The variability of  
498 wind resources in complex terrain and its relationship with atmospheric stability, *Energy Convers. Manag.*,  
499 222(July), 113249, doi:10.1016/j.enconman.2020.113249, 2020.

500 Radünz, W. C., Sakagami, Y., Haas, R., Petry, A. P., Passos, J. C., Miqueletti, M. and Dias, E.: Influence of  
501 atmospheric stability on wind farm performance in complex terrain, *Appl. Energy*, 282(PA), 116149,  
502 doi:10.1016/j.apenergy.2020.116149, 2021.

503 Richiardone, R., Giampiccolo, E., Ferrarese, S. and Manfrin, M.: Detection of Flow Distortions and Systematic  
504 Errors in Sonic Anemometry Using the Planar Fit Method, *Boundary-Layer Meteorol.*, 128(2), 277–302,  
505 doi:10.1007/s10546-008-9283-0, 2008.

506 Ruissi, R. and Bossanyi, E.: Engineering models for turbine wake velocity deficit and wake deflection. A new  
507 proposed approach for onshore and offshore applications, *J. Phys. Conf. Ser.*, 1222(1), doi:10.1088/1742-  
508 6596/1222/1/012004, 2019.

509 Santos, P., Mann, J., Vasiljević, N., Cantero, E., Sanz Rodrigo, J., Borbón, F., Martínez-Villagrasa, D., Martínez,  
510 B. and Cuxart, J.: The Alaiz experiment: untangling multi-scale stratified flows over complex terrain, *Wind Energy*  
511 *Sci.*, 5(4), 1793–1810, doi:10.5194/wes-5-1793-2020, 2020.

512 Sanz Rodrigo, J.: Flux-profile characterization of the offshore ABL for the parameterization of CFD models, *EWEA*  
513 offshore, 2011.

514 Sanz Rodrigo, J., Borbón Guillén, F., Gómez Arranz, P., Courtney, M. S., Wagner, R. and Dupont, E.: Multi-site

515 testing and evaluation of remote sensing instruments for wind energy applications, *Renew. Energy*, 53, 200–210,  
516 doi:10.1016/j.renene.2012.11.020, 2013.

517 Sanz Rodrigo, J., Cantero, E., García, B., Borbón, F., Irigoyen, U., Lozano, S., Fernande, P. M. and Chávez, R. A.:  
518 Atmospheric stability assessment for the characterization of offshore wind conditions, *J. Phys. Conf. Ser.*, 625,  
519 012044, doi:10.1088/1742-6596/625/1/012044, 2015.

520 Sathe, A., Gryning, S.-E. and Peña, A.: Comparison of the atmospheric stability and wind profiles at two wind farm  
521 sites over a long marine fetch in the North Sea, *Wind Energy*, 14(6), 767–780, doi:10.1002/we.456, 2011.

522 Sathe, A., Mann, J., Barlas, T., Bierbooms, W. A. A. M. and van Bussel, G. J. W.: Influence of atmospheric stability  
523 on wind turbine loads, *Wind Energy*, 16(7), 1013–1032, doi:https://doi.org/10.1002/we.1528, 2013.

524 Schmidt, J., Chang, C.-Y., Dörenkämper, M., Salimi, M., Teichmann, T. and Stoevesandt, B.: The consideration of  
525 atmospheric stability within wind farm AEP calculations, *J. Phys. Conf. Ser.*, 749, 012002, doi:10.1088/1742-  
526 6596/749/1/012002, 2016.

527 Schotanus, P., Nieuwstadt, F. T. M. and Bruin, H. A. R.: Temperature measurement with a sonic anemometer and its  
528 application to heat and moisture fluxes, *Boundary-Layer Meteorol.*, 26(1), 81–93, doi:10.1007/BF00164332, 1983.

529 Sorbjan, Z. and Grachev, A. A.: An Evaluation of the Flux--Gradient Relationship in the Stable Boundary Layer,  
530 *Boundary-Layer Meteorol.*, 135(3), 385–405, doi:10.1007/s10546-010-9482-3, 2010.

531 Stiperski, I. and Rotach, M. W.: On the Measurement of Turbulence Over Complex Mountainous Terrain,  
532 *Boundary-Layer Meteorol.*, 159(1), 97–121, doi:10.1007/s10546-015-0103-z, 2015.

533 Stull, R. B.: *An Introduction to Boundary Layer Meteorology*, Springer Netherlands., 1989.

534 Wilczak, J. M., Oncley, S. P. and Stage, S. A.: Sonic Anemometer Tilt Correction Algorithms, *Boundary-Layer*  
535 *Meteorol.*, 99(1), 127–150, doi:10.1023/A:1018966204465, 2001.

536 Zhan, L., Letizia, S. and Valerio Iungo, G.: LiDAR measurements for an onshore wind farm: Wake variability for  
537 different incoming wind speeds and atmospheric stability regimes, *Wind Energy*, 23(3), 501–527,  
538 doi:https://doi.org/10.1002/we.2430, 2020.

539