

Enhancing turbulent fluctuation measurement with tailored wind lidar profilers

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Abstract. This study evaluates the impact of an enhanced sampling rate on turbulence measurements using the Vaisala WindCube v2.1 lidar profiler. A prototype configuration, sampling four times faster than the commercial setup, was compared to the commercial WindCube v2.1 with reference measurements provided by a 2D sonic anemometer mounted on a measurement mast. Over the 47-day experiment, the prototype configuration showed performance similar to the commercial setup for key performance indicators (KPIs) like slope and coefficient of determination of mean wind speed compared to reference measurements, with both configurations meeting "best practice" threshold. However, for mean wind speed differences, the commercial configuration met the "best practice" level, while the prototype met the "minimum acceptance" criterion. Additionally, the data availability of the prototype configuration was 0.5% lower than that of the commercial configuration. Moreover, the increased sampling rate in the prototype lidar resulted in higher mean variance in instrumental noise compared to the commercial configuration. Despite this limitation, the mean noise-corrected along-wind variance measured by the prototype lidar was approximately 7% higher than that of the commercial lidar, **suggesting that the prototype might be better at capturing additional turbulent energy by resolving smaller eddies**. This effect was especially evident at higher wind speeds. Error metrics for the noise-corrected along-wind standard deviation in the prototype lidar were approximately 25% lower than those of the commercial configuration. However, the observed improvements of the prototype configuration in measuring turbulence fell short of expectations due to inherent limitations in the measurement process within the probe, where spatial and temporal filtering effects constrain the detection of turbulence at certain scales.

1 Introduction

Accurate turbulence data enable better understanding and control of wind flow patterns, optimizing the design, operation, and maintenance of wind energy systems. Furthermore, turbulence measurement plays a pivotal role in addressing key challenges within the wind energy sector. It aids in enhancing the efficiency and safety of wind turbine operations, minimizing wear and tear on vital components, and extending the lifespan of these costly assets. Additionally, improved turbulence measurement can facilitate more precise wind resource assessments, aiding in site selection and the overall planning of wind energy projects.

In the wind energy sector, the utilization of wind lidar profiler technology has gained significant traction in recent years, supplanting the traditional meteorological mast equipped with in-situ sensors like cup or sonic anemometers as the standard

25 means of measuring key mean wind properties, such as speed and direction. Wind lidar profilers present compelling advantages, including the potential for cost reduction compared to meteorological masts and the capacity to measure at similar or even greater heights above the ground.

Wind lidars profilers can be categorized according to their emission waveform, i.e., pulsed or continuous, and measuring technique, i.e., Doppler beam swinging (DBS) (Strauch et al., 1984) or velocity-azimuth display (VAD) (Browning and
30 Wexler, 1968). Measurement methods used by wind lidar profilers are fundamentally different from those used by cup or sonic anemometers. Anemometers estimate wind speed over a small volume of just a few cubic centimeters, whereas pulsed lidar profilers provide an average over a cylindrical probe several dozen meters long with a cross-sectional diameter of less than 1 cm (Fig. 1).

However, wind lidar profilers have yet to garner widespread acceptance for turbulence measurement, which remains a focal
35 point of ongoing research. In contrast to turbulence data derived from reference instruments such as sonic anemometers, turbulence data from lidar profiler measurements suffer from systematic errors induced by (i) the inter-beam effect, also known as the cross-contamination effect, (ii) the intra-beam effect, i.e., the averaging effect of the probe volume (Fig. 1) and, (iii), instrumental noise.

The inter-beam effect can result in either underestimation or overestimation of turbulence metrics, arising from the modulation
40 of energy associated with eddies of specific wavenumbers (Kelberlau and Mann, 2020). [This effect is particularly relevant in](#)

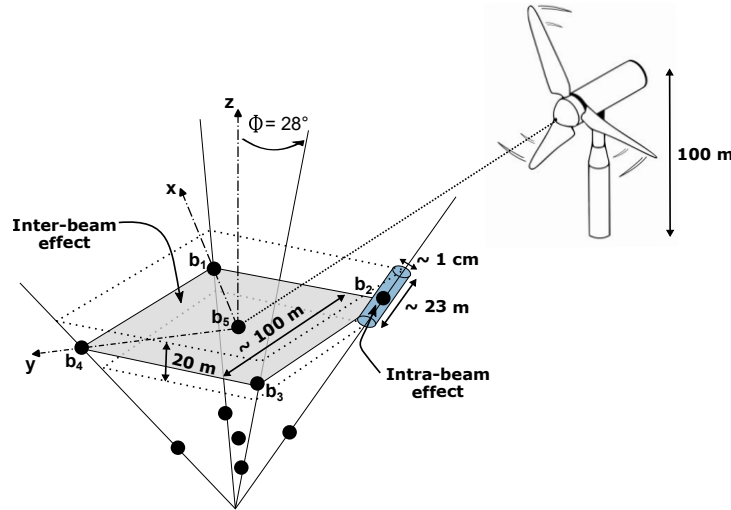


Figure 1. A schematic illustration of inter- and intra-beam effects in the WindCube v2.1 lidar profiler measurement process. The blue cylinder represents the probe volume, corresponding to the dimensions of the commercial lidar configuration. The positions of the five beams are labeled as b_i , where i ranges from 1 to 5. The inclination of the diverging beams (from beam 1 to beam 4) with respect to the vertical z -axis is $\phi = 28^\circ$. Beam 5 is aligned with the z -axis, while beams 1 and 3 are aligned with the x -axis, and beams 2 and 4 are aligned with the y -axis in the coordinate system of the instrument, as stipulated by the manufacturer.

the context of the assumption of instantaneous homogeneity, which underlies multi-beam lidar measurement techniques. Under this assumption, the turbulent field is considered spatially homogeneous across the beams at each instant in time, a condition that, if violated, can lead to inter-beam contamination. Any phase difference between the horizontal and vertical components

of an eddy significantly impacts the filtering of flow structures, potentially leading to both amplification or attenuation of their measured turbulent energy (Theriault, 1986; Gargett et al., 2009).

The intra-beam effect ~~generates~~ refers to a probe-time averaging phenomenon occurring within the lidar probe, leading to an underestimation of turbulence metrics. It arises from two anisotropic filtering processes: (1) spatial filtering due to averaging over the probe volume and (2) temporal filtering caused by averaging over the beam's pulse accumulation time, Δt , at a given measurement position. These two effects give rise to a transfer function, H , applied by the instrument on the signal measured within the probe. The transfer function includes a part due to time-averaging (the sinc term) and a part due to space-averaging (the Gaussian term), such that ~~(e.g., Kristensen et al., 2014):~~

$$|H|^2(\mathbf{k}) = \text{sinc}^2\left(\frac{\Delta t}{2}\mathbf{k} \cdot \mathbf{U}\right) \exp\left(-\left[\sigma_l^2(\mathbf{k} \cdot \mathbf{b})^2 + \sigma_r^2(\|\mathbf{k}\|^2 - (\mathbf{k} \cdot \mathbf{b})^2)\right]\right) \quad (1)$$

Here, \mathbf{k} is the turbulent structure wavevector, \mathbf{b} is the beam pointing vector, \mathbf{U} is the vector associated with the wind direction of magnitude U , and σ_l and σ_r represent the Gaussian weighting factors in the along-beam and cross-beam directions, respectively. A detailed mathematical derivation of Eq. 1 is provided in the supplementary material.

From Eq. 1, it follows that wind field structures with wavelengths smaller than σ_l in the along-beam direction are attenuated, as are those with wavelengths smaller than σ_r in the cross-beam direction. However, in the latter case, these structures are so small that the filtering effect becomes negligible, as the cross-section of the probe is approximately 1 cm (Fig. 1). Ultimately, assuming the Taylor frozen turbulence hypothesis, the wavevector domain that passes through the filter is defined by the intersection of two slices: one perpendicular to the direction of \mathbf{U} , which preserves structures longer than $\pi\Delta tU$, and another perpendicular to the direction of \mathbf{b} , which retains structures longer than σ_l . All other structures are filtered out.

Pulsed lidar profilers require several seconds to complete a full scanning cycle resulting in a low sampling rate that causes discrepancies between turbulence measurements taken by anemometers and those by lidar profilers (e.g., Peña et al., 2009). While the sampling rate governs how quickly the lidar progresses through a scan cycle, it is directly influenced by pulse accumulation time. Consequently, even if the sampling rate is increased, pulse accumulation can still limit the ability of the lidar to resolve small-scale turbulent structures. Since turbulent motion scales vary from milliseconds to hours and from centimeters to kilometers (e.g., Stull, 2000), it is crucial to account for both temporal and spatial filtering effects when assessing lidar-based turbulence measurements.

The concept of measuring turbulence using remote sensing instruments has gradually evolved since the early works in radar meteorology by Lhermitte (1962) and Browning and Wexler (1968). Lhermitte (1969) was the first to propose a method for inferring turbulence by analyzing the variance of radial velocity measurements through VAD scanning. Following this, Wilson (1970) conducted pioneering experiments using a pulsed Doppler radar to detect turbulence within the convective boundary

layer (0.1-1.3 km). However, these early measurements were limited to turbulence scales larger than the pulse volume and smaller than the scanning circle, and no validation against reference instruments was performed, questioning their reliability.

Kropfli (1986) expanded Wilson's approach to capture turbulence scales larger than the scanning circle by integrating data from multiple scans. Although initially developed for Doppler radar, these methods were later adapted for Doppler lidar. Eberhard et al. (1989) were the first to apply Wilson's and Kropfli's methods using lidar, and Gal-Chen et al. (1992) further refined the technique with a different scanning configuration. Despite these advancements, the significant probe length (around 100 m) limited studies to the convective boundary layer due to considerable probe volume averaging, especially near the ground. Frehlich (1994) and Frehlich (1994) demonstrated the averaging effect in the measurement of the structure function, showing that this effect becomes more pronounced at smaller separation distances. To address this limitation, research shifted ~~towards toward~~ understanding and mitigating probe volume averaging effects (e.g., Smalikho et al., 2005; Mann et al., 2010; Branlard et al., 2010). ~~Nowadays, modern lidarsystems have.~~ Smalikho et al. (2005) provided explicit formulae to account for the small-scale filtering effect of the finite probe volume in continuous-wave lidar, proposing three different methods for a staring lidar: (1) using the width of the Doppler spectrum, (2) the velocity structure function, and (3) the one-dimensional velocity spectrum. The expression for the structure function was derived under the assumption of local isotropy in the inertial subrange. Kristensen et al. (2011) later re-derived this expression, assuming a Lorentzian probe volume weighting function. Modern lidar systems have since reduced probe lengths to about approximately 30 m, but averaging effects still. However, averaging effects continue to pose challenges for turbulence measurements in the surface layer, where wind turbines operate (e.g., Mann et al., 2009; Sjöholm et al., 2009; Sathiyamoorthy et al., 2010). ~~These effects lead to an underestimation of variance derived from wind lidar compared to reference turbulence measurements, as demonstrated in (e.g., Mann et al., 2009; Sjöholm et al., 2009).~~

~~The present paper delves into the specific advancements pertaining to~~ This paper explores advancements in the Vaisala WindCube v2.1 lidar profiler. ~~A key modification is explored: an augmentation of,~~ focusing on a key modification: increasing the sampling rate, ~~achieved~~ by reducing the pulse accumulation time. This ~~customization is evaluated~~ enhancement is assessed for its impact on ~~the measurement of~~ measuring mean wind speed, data availability, and along-wind variance and its square root, i.e., the standard deviation. The ~~impact~~ latter is particularly important, as it is used in the wind power industry to compute turbulence intensity (TI), a critical metric for turbine load assessment, site suitability, and energy yield predictions. Additionally, the influence of instrumental noise on this modification is also evaluated to demonstrate is analyzed to confirm that the potential improvement improvements in turbulence estimates with the ~~version featuring increased sampling rate is not~~ due to higher sampling rate are not simply a result of increased noise.

The paper begins with a detailed overview of the data and methods, including the prototype configuration of the WindCube v2.1 and the field measurement setup. The study compares the prototype configuration against the commercial WindCube v2.1 and a sonic anemometer installed on a meteorological mast, which serves as the reference measurement. The methodology section then focuses on velocity spectra analysis, instrumental noise evaluation, and variance computation in instrument coordinates. Additionally, key performance indicators and error statistics used for validation are outlined. The results section presents findings on mean wind speed, data availability, standard deviation, variance and instrumental noise contributions. This is followed by a discussion of the implications, addressing both the advantages and challenges of a higher sampling rate.

The paper concludes with key takeaways on how the increased sampling rate enhances turbulence detection while considering measurement limitations and filtering effects.

110 2 Data and methods

2.1 Prototype configuration with increased sampling rate

The WindCube v2.1 lidar is designed for general atmospheric measurements, such as mean wind speed and direction, requiring a careful balance between temporal resolution, spatial resolution, and carrier-to-noise ratio (CNR). Its default sampling rate is optimized to ensure high data quality and availability across varying altitudes and atmospheric conditions while maintaining
115 system efficiency and manageable data processing.

The WindCube v2.1 employs the Doppler Beam Swinging (DBS) technique to measure wind speed. This method utilizes an optical switch that sequentially directs the lidar beam toward four ~~cardinal~~ directions (0° , 90° , 180° , and 270° ~~from True North~~ relative to the reference x -axis), each inclined at $\phi = 28^\circ$ from the vertical. A fifth beam is directed vertically upwards, resulting in wind measurements at five distinct positions (Fig. 1-2).

120 In its standard commercial configuration, the WindCube lidar collects data at each position for approximately $\Delta t = 0.8$ ~~seconds~~ s before switching to the next. Including transition times, a complete DBS scan is performed in 4 ~~seconds~~ s, yielding a line-of-sight (LOS) velocity sampling rate of 0.25 Hz (Table 1). This sampling rate is well-suited for capturing turbulent structures larger than 100 ~~meters~~ m. However, wind turbine components experience loads from turbulence across a wide range

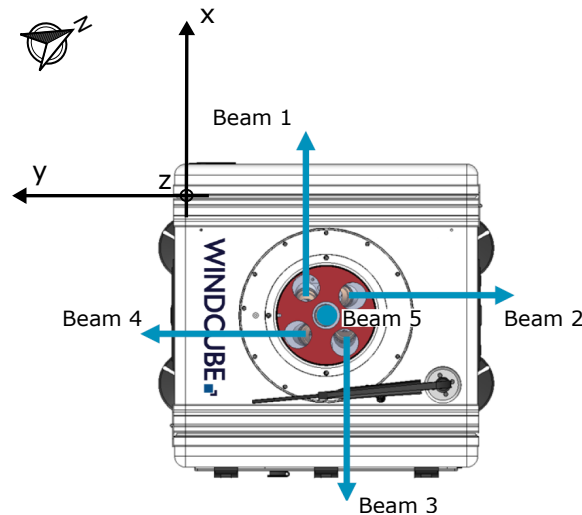


Figure 2. Top view of a WindCube v2.1 lidar showing the positions of its five beams. The x -axis is oriented from beam 3 towards beam 1, the y -axis extends from beam 4 towards beam 2, and the vertical z -axis points upward along beam 5. The arrow indicates North. For the present study, the primary x -axis of the lidars was oriented at -62° relative to North.

Table 1. LOS velocity measurement parameters for the commercial and prototype WindCube v2.1 configurations.

<u>Configuration</u>	<u>LOS Sampling Rate (Hz)</u>	<u>Accumulation Time (s)</u>	<u>LOS Samples per 30 min</u>	<u>Probe Length (m)</u>
<u>Commercial</u>	<u>0.25</u>	<u>0.8</u>	<u>450</u>	<u>23</u>
<u>Prototype</u>	<u>1.00</u>	<u>0.2</u>	<u>1800</u>	<u>23</u>

of scales. Increasing the sampling rate is crucial for broadening the velocity spectrum captured by the lidar, potentially enabling
125 the detection of ~~smaller-scale turbulence~~ extra turbulent energy that influences turbine performance.

Theoretically, a higher sampling rate improves temporal resolution and extends the resolved turbulence frequency range. However, for wind lidar profiler technology, this enhancement comes with trade-offs. The duty cycle, which represents the proportion of time the lidar transmits pulses, decreases as sampling rate increases, potentially reducing signal strength. Moreover, increasing the sampling rate requires a reduction in accumulation time, resulting in fewer pulses per sample and increasing
130 noise. The ~~default commercial~~ WindCube v2.1 configuration balances these factors to maximize data reliability. It integrates a high number of pulses per measurement to enhance signal quality, making it well-suited for general wind resource assessment. However, its probe length of approximately $L_{\text{probe}} = 23$ ~~meters-m~~ (Fig. 1, Table 1) limits its ability to resolve small eddies compared to point sensors like sonic anemometers.

In response to the demand for capturing ~~the energy of smaller eddies~~ additional turbulent energy, we developed a modified
135 version of the WindCube v2.1 that operates four times faster, achieving a LOS velocity sampling rate of 1 Hz. This modification was achieved by reducing the accumulation time for data collection from each beam in conjunction with a reduction in the number of transmitted pulses. The factor of 4 was chosen as a compromise between increasing temporal resolution and maintaining an acceptable CNR and data availability. This choice is intended to keep wind measurements comparable to those from the commercial configuration while enabling the ~~investigation of smaller-scale turbulence~~ capture of additional turbulent
140 energy. The actual impact on measurement performance will be assessed in the study.

2.2 Field measurement

2.2.1 ~~Measurement site and data collection~~

The field measurement campaign was ~~carried out by DNV-GL~~ conducted by DNV at the lidar validation test site in Janneby, Germany (Fig. 3). ~~Due to its flat terrain, the site features orography-undisturbed flow, which is suitable for the verification~~
145 ~~trialsof lidar systems. The site has relatively~~ The site's flat terrain ensures orography-free flow, making it ideal for lidar verification trials. It offers good exposure to largely undisturbed wind conditions ~~, i.e., undisturbed winds from almost all sectors. The elevation of the site is from most directions. Situated just~~ a few meters above mean sea level, ~~and the surface roughness is low due to the primarily~~ the site features low surface roughness due to its predominantly agricultural land use (Fig. 3a). Two wind turbines (WT N100 and WT N117~~in~~; Fig. 3a) are ~~installed in the vicinity of the~~ located near the
150 meteorological mast. The closest wind turbine is located Their wake-affected wind sectors are shown in blue in Fig. 4 and lie outside the sectors selected for turbulence analysis (gray areas in Fig. 4, Section 2.5). The closest turbine is 210 m from the

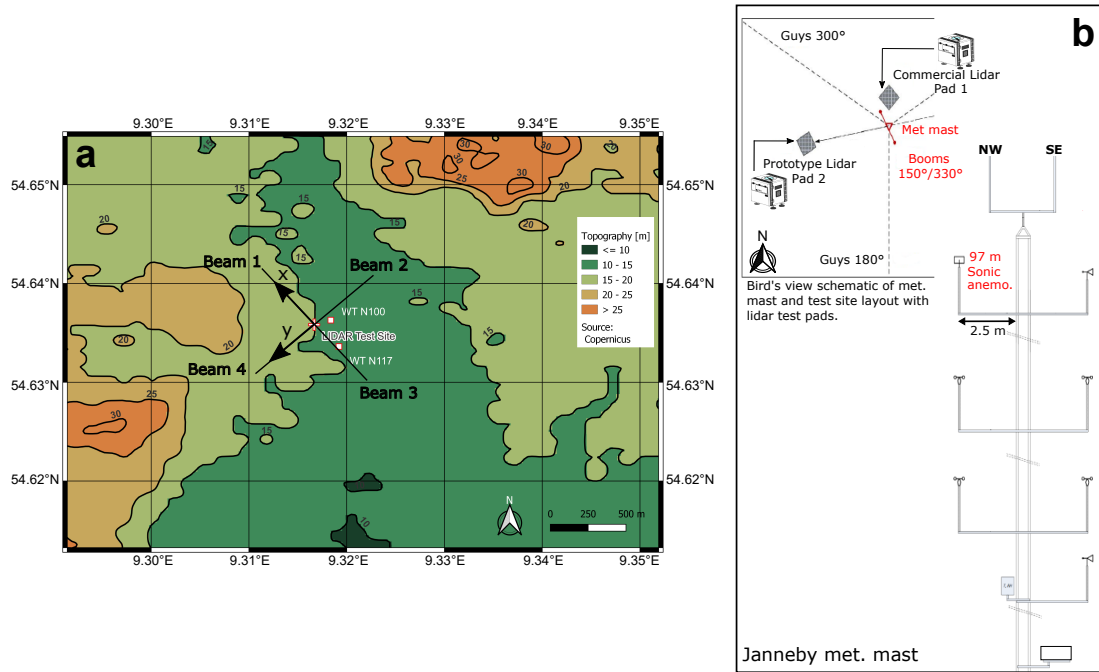


Figure 3. a: Test site location at Janneby, Germany. Black lines indicate the beam orientations for the commercial and prototype configurations. The x and y axes of the instrument coordinate system (see Fig. 2) are marked with black arrows. b: Configuration of the meteorological mast, showing the position of the sonic anemometer. NW and SE denote the north-west and south-east directions. The schematic in panel b also provides a bird's-eye view of the meteorological mast and test site layout, including the lidar test pads.

mast. A few ~~human constructions~~ small human-made structures (e.g., houses, sheds), ~~not exceeding all under~~ 15 m in height, are ~~located approximately~~ situated about 500 m southwest of the mast.

The meteorological mast is a 100 m, 3-fold guyed lattice tower with a constant face width of 0.4 m. It is equipped with
 155 six MEASNET-calibrated Thies First Class Advanced cup anemometers (No. 4.3352) and a Thies 2D sonic anemometer (No. 4.3830). However, only the Thies 2D sonic anemometer is used in this study to provide reference measurements of mean wind speed and turbulence, as the cup anemometers data are not available. The mounting arrangements are consistent with the currently valid IEC and IEA recommendations for the use of anemometry at meteorological masts. As shown in Fig. 3b, the sonic anemometer is pointing towards 150° from True North and is mounted at 97 m above ground, which corresponds to the
 160 average hub height of modern land-based wind turbines. The sonic anemometer was set to record continuous horizontal wind speed and direction at sampling rates of 4 Hz.

Adjacent to the measurement mast, both the commercial lidar configuration, and a prototype version with an enhanced sampling rate were installed 3 m and 13 m apart the mast respectively. ~~The prototype configuration was set to record the LOS velocity four times faster than the commercial configuration, as described in Section 2.1. The lidar was~~ Both lidars were aligned
 165 such that beams 1 and 3, which correspond to the x -axis (Fig. 2), were oriented at -62° from True North (Fig. 3). According to

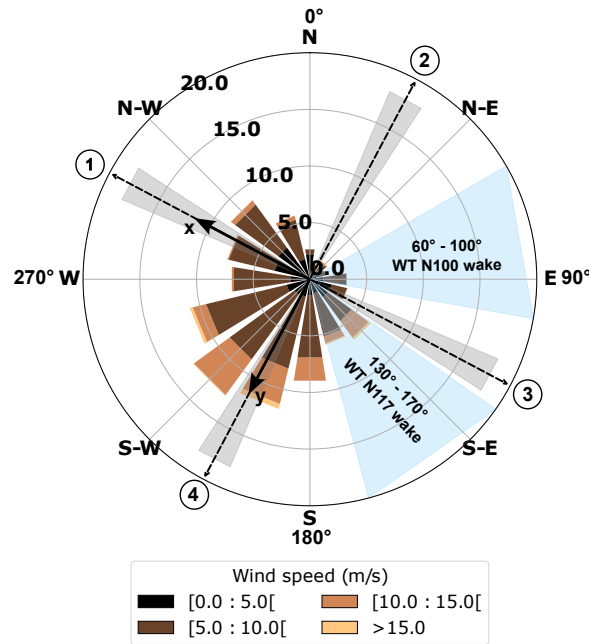


Figure 4. Wind rose showing wind data recorded over 47 days by the sonic anemometer at 97 m above ground level. Gray shaded areas indicate the wind sectors selected for the turbulence analysis in this study, corresponding to events when the wind was aligned ($\pm 5^\circ$) with either the beam pair 1-3 (aligned with the x -axis) or beam pair 2-4 (aligned with the y -axis), as numbered in circles. Blue shaded areas indicate wind sectors contaminated by the nearby wind turbines WT N100 and WT N117 (Fig. 3a).

the manufacturer’s recommendation, the x -axis is the primary axis and should be oriented relative to North. Beams 2 and 4 are fixed along the y -axis.

The field measurement campaign was conducted over two periods: from 12 to 25 November 2021, and from 07 December 2021 to 10 January 2022. These two measurement periods were combined to form a 47-day dataset. To facilitate a comparison
170 of turbulence measurements, the sonic-derived wind dataset was resampled to match the sampling rate of the LOS velocities measured by the prototype configuration. This ensures that similar turbulence time scales are captured when calculating and comparing turbulence estimates. Therefore, the sonic anemometer measurements were resampled at 1 Hz.

The 47-day dataset was ~~divided into 2256 subsets of 30-minute data records, with each subset containing segmented into 2,256 subsets, each comprising 30-min of data. For each subset, the commercial lidar provided 450 and 1800 measurement points for~~
175 ~~wind data acquired at~~ measurement points, while the prototype lidar provided 1,800 points, corresponding to their respective sampling rates of 0.25 Hz and 1 Hz ~~, respectively. The choice of a 30-minute window deviating from the standard 10-minute window typically (see Table 1). The selection of a 30-min window, rather than the standard 10-min interval commonly used in the wind energy industry was informed by considerations of reduction of, was guided by the aim of reducing random errors in turbulence measurements, as discussed by following the recommendations of~~ Lenschow et al. (1994).

180 2.3 Velocity spectra

Power spectral density of the velocity, i.e., the velocity spectra, provide valuable information about the distribution of turbulent kinetic energy across different scales of motion within the wind flow. This understanding helps in characterizing turbulence and its effects on wind turbine performance and structural loads.

Velocity spectra were computed using Welch's method (Welch, 1967). This method computes an estimate of the spectrum
185 by dividing the data into overlapping segments, computing a modified periodogram for each segment and averaging the periodograms. The Hann window with 50% overlap was applied to each segment to reduce spectral leakage and improve frequency resolution. The 50% overlap is a reasonable trade off between accurately estimating the signal power, while not over counting any of the data.

~~Following the recommendations of Kelberlau and Mann (2020), Fitting~~ turbulence velocity spectra ~~computed from the~~
190 ~~lidar-derived reconstructed velocity component should not be fitted~~ derived from lidar-reconstructed velocity components
to turbulence models ~~;-This is~~ should be avoided due to the inter-beam effect, which distorts the spectra ~~;-rendering them~~
~~physically meaningless~~ and complicates their physical interpretation. Therefore, such spectra were not considered in this study.
The focus was on velocity spectra $S_i(f)$ derived from the LOS velocities measured by beam i . The primary limitation in this
approach is the intra-beam effect. Spectra were computed for each ~~30-minute~~ 30-min subset of data.

195 The spectra, $S_i(f)$, were fitted by a parametric expression (Teunissen, 1980; Olesen et al., 1984; Tieleman, 1995) in the frequency domain f , to which we add a component N_i associated with the power spectral density of instrumental noise of the LOS velocity measured by beam i (see section 2.4):

$$S_i(f) = \frac{m}{(1 + nf)^\beta} + N_i \quad (2)$$

The coefficient m primarily controls the vertical scaling or amplitude of the spectrum whereas n influences the rate at which
200 the function decays as f increases. The exponent β determined the shape of the spectrum.

Three different weighting schemes were considered: an unweighted scheme, a low-frequency weighted scheme with weights proportional to the logarithm of the frequency, and a high-frequency weighted scheme with weights inversely proportional to the logarithm of the frequency. Assessing the fitting accuracy included comparing the variance obtained from the integrated fitted spectra with the measured spectra, and calculating their absolute relative differences.

205 2.4 Instrumental noise

Lidar measurements are inherently influenced by signal noise and potential variations in aerosol fall speeds, both of which contribute additional terms to the observed variance. Assuming that all atmospheric flow contributions to the observed LOS velocity variance within the considered short timescales are of a turbulent nature, the variance $\sigma_{b_i}^2$ of the LOS velocity measured by beam i , can be expressed as the sum of three independent terms (Doviak and Zrnic, 1993):

$$210 \quad \sigma_{b_i}^2 = \sigma_{p_i}^2 + \sigma_{n_i}^2 + \sigma_{d_i}^2 \quad (3)$$

Here, $\sigma_{p_i}^2$ represents the net contribution from atmospheric turbulence at scales measurable by the lidar (Brugger et al., 2016), $\sigma_{n_i}^2$ denotes the variance associated with instrumental noise, and $\sigma_{d_i}^2$ accounts for the variance caused by variations in aerosol terminal fall speeds within the measurement-probe volume. However, $\sigma_{d_i}^2$ can typically be neglected, as particle fall speeds are generally less than 1 cm/s (e.g., Bodini et al., 2018). Noise will-be-has-been identified through two different methods: a spectral
215 approach and an autocorrelation approach, as accurately identifying the variance of noise is critical to our study.

2.4.1 Spectral method

Instrumental noise is a critical factor in the spectral analysis of velocity time series. In the spectrum of a velocity time series, this noise typically manifests as a flattening of the spectrum at higher frequencies, indicating a white noise characteristic that contributes equally across these frequencies (e.g., Thomson et al., 2012; Durgesh et al., 2014; Guerra and Thomson, 2017;
220 McMillan and Hay, 2017; Thiébaud et al., 2020). At lower frequencies, the spectrum is usually dominated by the actual signal, which may show a characteristic decay or specific features related to the physical process being measured, such as turbulence. As frequency increases, the influence of the instrumental noise becomes more prominent, leading to a flattened spectral region where the noise dominates.

In Eq. 2, N_i represents the constant power spectral density of noise, which contributes to the spectral flattening observed
225 at higher frequencies. The variance of the noise depends on the technical characteristics of the device measuring the velocity, such as Nyquist velocity, the signal spectral width, the number of pulses and points per range gate, and the signal-to-noise ratio. Theoretical expressions for the variance of this noise can be derived and subsequently removed from the computed turbulence metrics to improve accuracy (Pearson et al., 2009; O'Connor et al., 2010; Bodini et al., 2018, 2019; Wildmann et al., 2019). However, the technical specifications of lidar profilers are no longer openly shared with users, making it impossible to
230 evaluate this noise theoretically. To address this, it is essential to evaluate the noise using an alternative method, such as the spectral approach employed in this study. This approach is comparable to the method proposed by (e.g. Richard et al., 2013; Durgesh et al., 2014). It enables the determination of the power spectral density of noise, N_i , associated with the LOS velocity measured by beam i . Subsequently, the variance of the instrumental noise, $\sigma_{n_i}^2$, can be derived by multiplying N_i by the Nyquist frequency, f_N , such as (e.g., McMillan and Hay, 2017):

$$235 \quad \sigma_{n_i}^2 = N_i f_N \quad (4)$$

2.4.2 Autocorrelation function method

An alternative method for computing the variance of the instrumental noise involves the calculation of the auto-correlation function (ACF) of the squared LOS velocity time series, as proposed by Lenschow et al. (2000). The ACF quantifies the similarity between a signal and its time-shifted versions across various time lags. This measure provides insight into how much
240 of the signal correlates with its past values, which is essential for distinguishing between the noise and signal components.

According to Lenschow et al. (2000), after calculating the ACF, the ACF values (excluding the first lag) are fitted to a 2/3 power-law function. This power-law model describes the decay of correlation over time, allowing for the extraction of a

coefficient that characterizes how the correlation diminishes as the time lag increases. From this power-law fit, the value of the ACF as the lag tends to zero is estimated by extrapolation of the fitted model. This value is associated with the signal variance.

245 Subsequently, the total variance of the signal is calculated. The instrumental noise variance, $\sigma_{n_i}^2$, is then determined by subtracting the signal variance, as derived from the fitted power-law model, from the total variance. This process enables the separation of the signal and noise contributions to the overall variance. ~~However, this method performs correctly only if the range in which the turbulent cascade occurs is fully captured. This condition is not guaranteed with wind lidar profiler measurements, as the intra-beam effect disturbs the inertial range of turbulence where the cascade takes place.~~ The application
250 of the ACF method requires that the data be stationary. To verify this assumption, a stationarity test was performed for each 30-min subset using the Augmented Dickey-Fuller (ADF) test. The null hypothesis of the ADF test states that the data are non-stationary which was assessed through the determination of a significance level, commonly set at 0.05 or 5% (e.g. Hayat, 2010).

2.5 Computation of the variance in instrument coordinates

255 The conventional method for computing variance and standard deviation (the square root of variance) from wind lidar profiler measurements relies on deriving second-order statistics from the reconstructed instantaneous velocity components based on LOS velocities. This approach inherently combines, at each time step, measurements taken at sampling points separated by several tens of meters, depending on the height level of interest. The assumption of instantaneous flow homogeneity (inter-beam effect) introduces an uncertainty in the derived statistics, which is difficult to quantify and can lead to either an overestimation
260 or underestimation of the standard deviation, depending on the frequency and flow configuration. Additionally, this traditional method is affected by both intra-beam filtering and instrumental noise. Crucially, because variance is computed from the reconstructed instantaneous velocity components, it does not account for the noise-induced variance present in the LOS velocity time series which will result in overestimation of variance.

The combined influence of the inter-beam effect, intra-beam effect, and instrumental noise can result in variance estimates
265 derived from the traditional approach that may appear to align more closely with those derived from a sonic anemometer, but for reasons unrelated to the actual turbulence characteristics. Consequently, the benefits of an increased sampling rate for turbulence measurement using a lidar profiler cannot be accurately assessed with this approach.

The variance method, as referred to in the studies (e.g., Stacey et al., 1999a, b; Lu and Lueck, 1999; Rippeth et al., 2002; Guerra and Thomson, 2017; Thiébaud et al., 2022), offers an alternative to the traditional approach for computing variance. This
270 method calculates the second-order statistics of the three velocity components by deriving them directly from the second-order statistics of the LOS velocities. Unlike the traditional approach, the variance method is unaffected by the inter-beam effect. However, it is still influenced by the intra-beam effect and instrumental noise. Notably, the impact of instrumental noise can be identified and removed. Hereafter, a hat notation is used to denote standard deviation or variance derived from this method.

The variance method enables the calculation of the variances, $\hat{\sigma}_x^2$ and $\hat{\sigma}_y^2$ of the velocity components u_x and u_y (in instrument
275 coordinates) as:

$$\hat{\sigma}_x^2 = \frac{1}{2 \sin^2 \phi} \left(\sigma_{p_3}^2 + \sigma_{p_1}^2 - 2 \cos^2 \phi \sigma_{p_5}^2 \right) \quad (5)$$

$$\hat{\sigma}_y^2 = \frac{1}{2 \sin^2 \phi} \left(\sigma_{p_2}^2 + \sigma_{p_4}^2 - 2 \cos^2 \phi \sigma_{p_5}^2 \right) \quad (6)$$

where $\sigma_{p_i}^2 = \sigma_{b_i}^2 - \sigma_{n_i}^2$ (Eq. 3), is the variance of the LOS velocity recorded by beam i , corrected for the variance of instrumental
280 noise.

In this paper, we restrict the application of the variance method to situations where the wind aligns ($\pm 5^\circ$) with a single
pair of opposite beams (either pair 1-3 or pair 2-4) of the lidar profilers. This alignment condition was met in 17.1% of the
cases ~~-(gray areas in Fig. 4).~~ Under these conditions, it can be reasonably assumed that the covariance term, $\hat{\sigma}_{uv}$ (where v
represents the cross-wind velocity), which corresponds to $\hat{\sigma}_{xy}$ in this specific condition, is negligible (e.g., Newman et al.,
285 2016). Specifically, when the wind aligns with beams 1 and 3, we have $\hat{\sigma}_u^2 = \hat{\sigma}_x^2$. Conversely, when the wind aligns with beams
2 and 4, it follows that $\hat{\sigma}_u^2 = \hat{\sigma}_y^2$. For brevity, we use $\hat{\sigma}^2$ in place of $\hat{\sigma}_u^2$ hereafter. The standard deviation, $\hat{\sigma}$, is then compared to
the along-wind standard deviation, σ , which is derived from sonic anemometer measurements.

2.6 Key performance indicators and acceptance criteria

The first step of our analysis focuses on key performance indicators (KPIs) ~~-, applied-related~~ to mean wind statistics, such
290 as wind speed, ~~that are the~~. These include mean differences, slope, ~~or-and~~ the coefficient of determination (R^2) at reference
heights corresponding to sonic anemometer measurements. ~~DNV-GL has defined~~ The verification process follows standard
lidar performance requirements set by DNV and NORSEWInD (2009), which define acceptance criteria (ACs) as either "best
practice" ~~and-or~~ "minimum allowable tolerances." ~~These criteria, applied~~ Applied to wind speed, these criteria flag any KPIs
~~outside the exceeding the defined~~ thresholds as "deviations". Table 2 summarizes the ACs established by ~~DNV-GL~~ DNV,
295 which are tested in this paper for ~~the wind speed KPI~~ wind speed KPIs.

Additionally, the paper addresses data availability. Data availability is defined as the ratio of valid data points returned by the
lidar to the maximum number of possible points that could be acquired during the test. To pass the test, ~~DNV-GL the standard~~
lidar performances set the data availability threshold at 90% (Table 2).

2.7 Error statistics metrics

300 This paper focuses on turbulence measurements, specifically the standard deviation ~~of wind-and~~ variance of the along-wind
velocity, obtained from both the commercial and prototype lidars. These measurements are compared to the standard deviation
provided by the reference instrument, ~~the~~ the sonic anemometer. To assess the accuracy and reliability of the lidar turbulence
measurements, various error statistics are used. These include ~~Root-Mean-Square-Error (RMSE), which~~:

305

- Root Mean Square Error (RMSE): quantifies the average magnitude of ~~errors~~; ~~Mean Absolute Error (MAE)~~, which the errors.

$$\text{RMSE} = \sqrt{\frac{1}{M} \sum_{i=1}^M (X_i - Y_i)^2} \quad (7)$$

- Mean Absolute Error (MAE): calculates the average absolute difference between predicted and observed values; ~~bias~~, which.

$$\text{MAE} = \frac{1}{M} \sum_{i=1}^M |X_i - Y_i| \quad (8)$$

310

- Bias: represents the systematic error between the lidar and reference measurements; ~~the coefficient of determination, R²~~, which.

$$\text{Bias} = \frac{1}{M} \sum_{i=1}^M (X_i - Y_i) \quad (9)$$

- Coefficient of determination (R²): indicates the proportion of variance in the lidar measurements explained by the reference data; ~~and relative error, which expresses the error as a percentage of the reference measurement.~~

315

$$R^2 = 1 - \frac{\sum_{i=1}^M (X_i - Y_i)^2}{\sum_{i=1}^n (Y_i - \bar{Y})^2} \quad (10)$$

Here, X_i represents the lidar measurement, Y_i the corresponding reference measurement from the sonic anemometer, \bar{Y} is the mean of the reference measurements, and M is the number of turbulence estimates. Together, these statistical metrics provide a comprehensive evaluation of the lidar's performance in capturing turbulence characteristics relative to the reference instrument.

Table 2. Acceptance criteria for KPI of mean wind speed in wind lidar profiler certification.

KPI - Wind speed	Definition	Best practice	Minimum	Deviation
Difference	Percentage difference in mean wind speeds between lidar and reference over the verification campaign, relative to the campaign mean wind speed.	< 1%	[+1.5] <u>[-1.5]</u> %	> 1.5%
Slope	Slope from single-variable regression, constrained to pass through the origin.	[0.98–1.02]	[0.97–1.03]	< 0.97 or > 1.03
R ²	Correlation coefficient from single-variable regression.	> 0.98	> 0.97	≤ 0.97
<u>Data availability</u>	<u>Mean percentage of available data points in each 30-min subset, relative to the total number of possible records.</u>	<u>≥ 90%</u>	~	<u>< 90%</u>

3.1 Mean wind speed and data availability

The first step in proposing enhancements to lidar technology is to evaluate their impact on mean wind speed measurements. Fig. 5a illustrates that the mean vertical wind speed profiles measured by both configurations are closely aligned. However, the difference between the mean wind speed measurements provided by the commercial configuration and the reference measurement (black cross in Fig. 5) at the reference altitude is smaller, amounting to 0.98%, compared to a 1.41% difference for the prototype configuration. These results demonstrate that the commercial configuration closely matches the "best practice" AC criterion for the difference in mean wind speed, while the prototype configuration, with a larger difference, only meets the "minimum" criterion (Table 2 and Table 3).

Moreover, the commercial configuration exhibits data availability ranging from 99.5% at the lowest measurement height, i.e., 40 m above the ground, to 93.0% at the highest, i.e., 200 m above the ground, with an overall vertical average availability of 98.2% (Fig. 5b). Similarly, the prototype configuration follows this trend, with data availability decreasing with altitude. The prototype achieves a vertical average availability of 97.7%, with a minimum of 92.3% recorded at the highest measurement altitude. The prototype configuration consistently shows data availability that is, on average, 0.5% lower than the commercial configuration at nearly all measurement altitudes. Both lidar configurations exceed the 90% data availability threshold set by DNV-GL DNV and NORSEWInD (2009).

Fig. 6 presents the linear regression of the 30-minute 30-min averaged wind speed measured by both lidar configurations in comparison to the reference instrument. Both the commercial and prototype configurations match the "best practice" criteria, with slope values of 1.0 and R² values of 0.9847 for the commercial configuration. The prototype configuration shows values that are 1% lower for the slope and almost-similar R², but these differences are minimal and still within the acceptable range for "best practice."

3.2 Impact of sampling rate on turbulence energy capture

The impact of increasing the sampling rate on turbulence measurement can initially be assessed using data from a sonic anemometer, specifically through the computation of along-wind velocity spectra. Integrating these spectra provides the along-wind variance, σ^2 . Fig. 7 illustrates the individual spectra and the mean spectrum averaged over the 47-day dataset in both

Table 3. Acceptance criteria for KPI achievement applied on mean wind speed associated with the commercial and prototype configurations: ✓✓ denotes "best practice" and ✓ indicates "minimum" acceptance, as defined in Table 2.

<u>Configuration</u>	Difference	Slope	R ²	Data availability
Commercial configuration-	✓✓	✓✓	✓✓	✓✓
Prototype configuration-	✓	✓✓	✓✓	✓✓

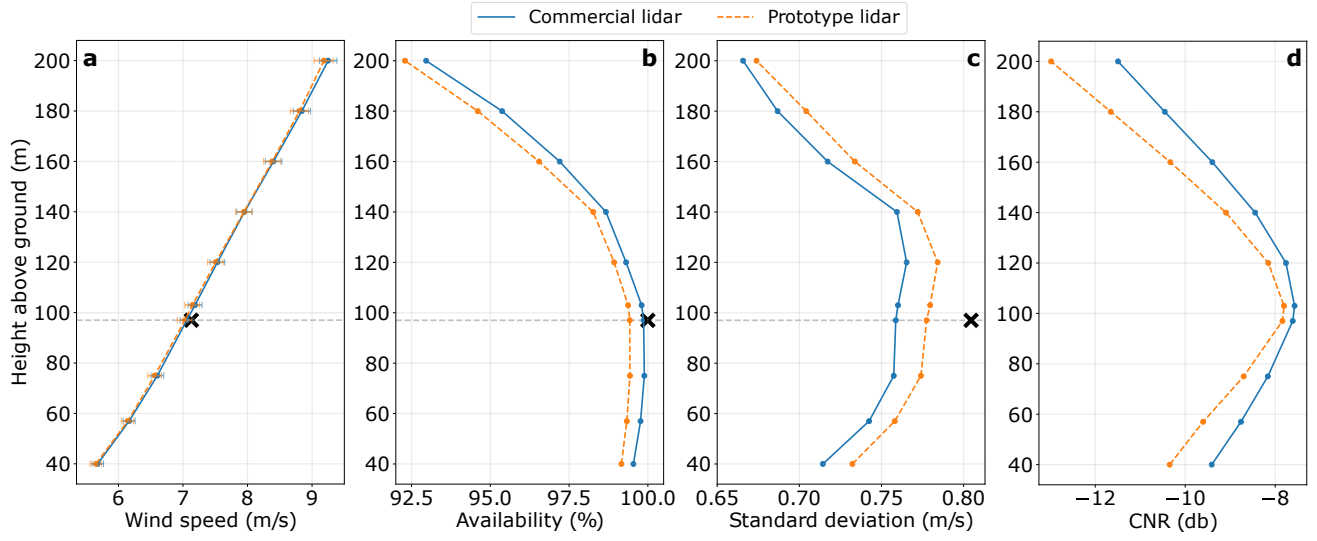


Figure 5. Mean vertical profiles, averaged across the 47-day dataset, of wind speed (a), data availability (b), ~~and~~ standard deviation derived from the variance method (c), and CNR (d) measured using the commercial (solid blue curves) and prototype (dashed orange curves) configurations. In (a), error bars represent 95% confidence intervals computed via bootstrapping, illustrating the statistical uncertainty of the mean. The black crosses represent the reference measurements from the sonic anemometer, and the grey dashed vertical line marks its position at 97 m above ground.

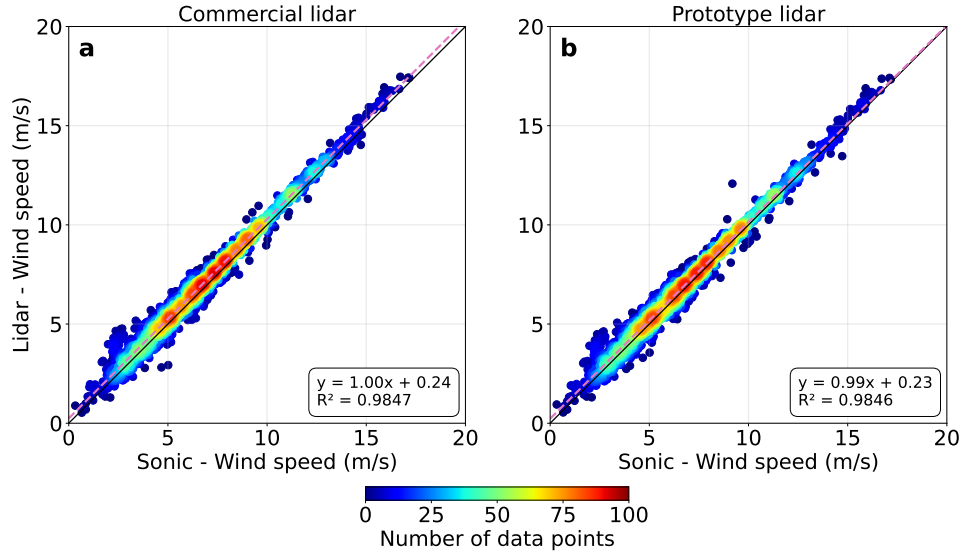


Figure 6. Scatter plots of the ~~30-minute~~ 30-min averaged wind speed measurements over the 47-day campaign, comparing the commercial lidar (a) and prototype lidar (b) with the reference sonic anemometer.

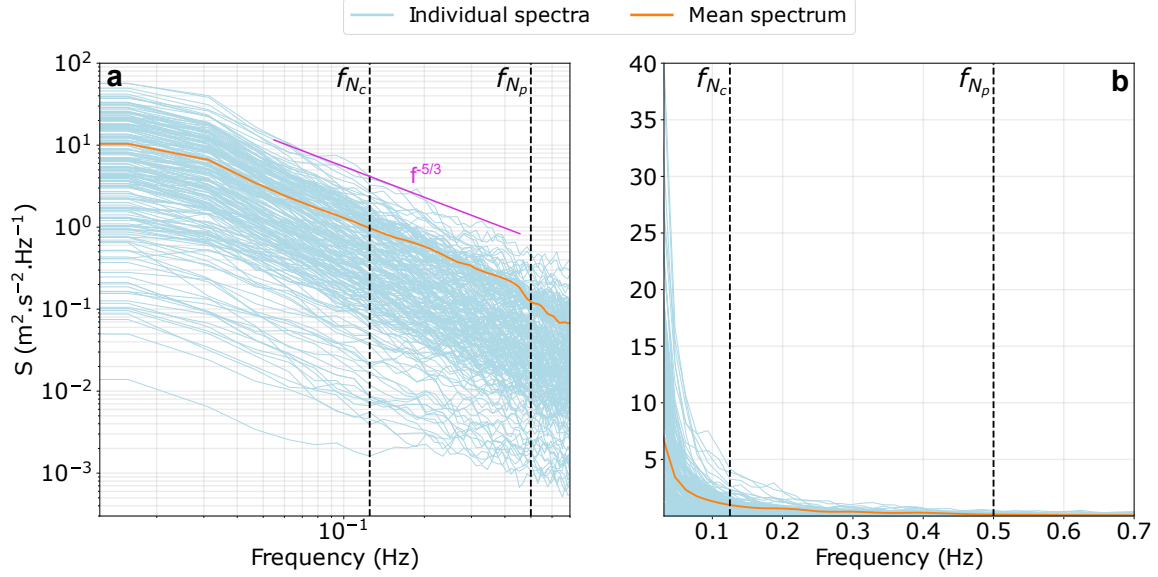


Figure 7. Individual spectra (light blue curves) and mean spectrum (orange curve) measured by the sonic anemometer over the 47-day measurement campaign, presented in log-log (a) and linear (b) formats. Vertical black dashed lines indicate the Nyquist frequencies, f_{N_c} and f_{N_p} , for the commercial and prototype lidar configurations respectively. The pink solid line in panel (a) shows the classic spectral slope $f^{-5/3}$.

log-log and linear formats. The mean spectrum clearly follows the $f^{-5/3}$ slope, confirming the presence of the energy cascade (Fig. 7a).

The linear representation (Fig. 7b) highlights that most of the energy, associated with larger eddies, is concentrated in the frequency range from 0 to $f_{N_c} = 0.125$ Hz, corresponding to the Nyquist frequency of the LOS velocity in the commercial lidar configuration. However, additional energy, associated with smaller eddies, exists within the range from f_{N_c} to $f_{N_p} = 0.5$ Hz, the latter being the Nyquist frequency of the prototype lidar configuration.

To quantify this effect, the variance was computed by integrating the spectra over two frequency ranges. First, the integration from 0 to f_{N_c} simulated the variance measurable by a sonic anemometer with a sampling rate equivalent to the commercial lidar. This yielded a mean variance of $0.47 \text{ m}^2/\text{s}^2$ $0.4712 \text{ m}^2/\text{s}^2$. Second, the integration from 0 to f_{N_p} simulated the variance measurable with a sampling rate equivalent to the prototype lidar, resulting in a mean variance of $0.63 \text{ m}^2/\text{s}^2$ $0.6314 \text{ m}^2/\text{s}^2$. This comparison indicates that increasing the sampling rate by a factor of 4, relative to the commercial lidar configuration, could capture an additional 34% of the ~~energy associated with smaller eddies~~ variance. However, this represents the maximum possible improvement, as it is derived from measurements using a sonic anemometer, which is not affected by technical limitations such as the ~~probe length~~ space-time volume averaging inside the probe of a wind lidar profiler.

3.3 LOS velocity spectra

The determination of the instrumental noise from the spectral method involves computational fitting of the LOS velocity spectra using a parametric expression (Eq. 2). Three weighting schemes were systematically explored to enhance fitting accuracy and minimize errors relative to the measured spectra. Fig. 8a illustrates an example of the three weighting scheme applied to a measured spectrum. This iterative process was conducted across both lidar configurations, yielding consistent results described hereafter.

The fitted spectra closely matched in the low-frequency domain, approximately up to $f = 0.1$ Hz, but strong divergences were observed thereafter. The low frequencies weighted scheme produced a curve substantially below the measured spectra at higher frequencies, whereas the unweighted scheme yielded a curve slightly above the measured spectra in this frequency range. In contrast, the high frequencies weighted scheme provided a fit that closely matched the measured spectra across all frequencies and exhibited the lowest mean error. For instance, when applied to the prototype lidar, the mean variance was $0.2321 \text{ m}^2/\text{s}^2$ for all integrated fitted spectra using the high-frequency weighted scheme, compared to $0.2262 \text{ m}^2/\text{s}^2$ for all integrated measured spectra. This results in an absolute error of 2.6%. Conversely, not employing any weighting during the fitting process resulted in an absolute error between the mean variance nearly three times higher, at 8.5%. Assigning weights to the low frequencies resulted in a mean absolute error exceeding six times that of the high-frequency weighted scheme, at 16.9%. Thus, the high-frequency weighted scheme was chosen for the fitting. An example of this fitting applied to individual

LOS velocity spectra for both the commercial and prototype configurations is shown in Fig. 8b. This weighted

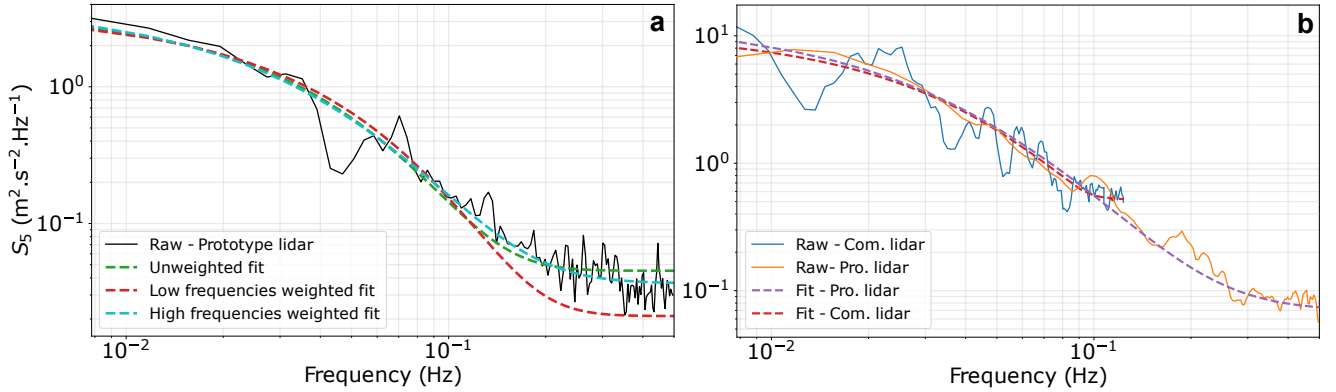


Figure 8. (a) Individual-LOS velocity spectrum (solid-black)-measured by beam 5 of the prototype lidar (solid black), fitted with using Eq. 2 with three weighted-different weighting schemes: unweighted fit (dashed green), low-frequencies-low-frequency weighted fit (dashed red), and high-frequencies-high-frequency weighted fit (dashed blue). This panel corresponds to the study focused on selecting the optimal weighting scheme. (b) Individual-The optimal scheme (high-frequency weighted) is applied to LOS velocity spectrum measured by beam 5 of the commercial lidar (blue curve) and the prototype lidar (orange curve) fitted with the high-frequencies-weighted scheme.

scheme enabled the systematic identification of the plateau at higher frequencies, characteristic of white noise. Other weighting schemes did not consistently exhibit this plateau, making it challenging to reliably determine the value of N_i .

3.4 Instrumental noise

3.4.1 Carrier-to-noise ratio

380 Fig. 5d shows that the CNR of the prototype lidar is consistently lower than that of the commercial system throughout the
altitudes of measurement, with a mean value across all heights of -8.9 dB. On average, the prototype lidar's CNR is 8.5% lower,
indicating a weaker signal and thus a higher relative noise level compared to the commercial lidar. Despite this difference in
magnitude, the mean vertical profiles of CNR for both lidars follow the same trend: the CNR reaches a minimum near the
ground, increases to a maximum at approximately 100 m above ground, and then decreases again, reaching another minimum
385 at the highest measurement altitude.

3.4.2 Comparison of the spectral and ACF methods

The spectral method yields a median variance that is 1.5 times higher than that of the ACF method for the commercial ~~lidar~~
~~and twice as high for the prototype lidar and prototype lidars~~ (Table 4). While this suggests differences in how each method
characterizes noise, the spectral method also results in a mean instrumental noise that is 30-40% lower than that of the ACF
390 method, indicating variations in the way noise is estimated. Moreover, the spread of mean values is notably narrower when
using the spectral method, particularly for the commercial lidar, where it is reduced by half compared to the ACF method. This
suggests a potential advantage in terms of consistency and stability. Given these observations, we used the spectral method to
correct the measured variance, as it appeared to provide more stable estimates of instrumental noise.

To validate the applicability of the ACF method and investigate the higher spread of mean values of instrumental noise
395 associated with this method, the ADF test was applied to each 30-min data subset. The results show that approximately 8%
of the subsets yielded p-values just above the 0.05 significance threshold, though none exceeded 0.06. This indicates that, for
these subsets, the null hypothesis of non-stationarity (see Section 2.4.2) could not be rejected. Consequently, they cannot be
confidently considered stationary, and the ACF method is not strictly valid for them. This limitation may partly account for
the higher variability in noise estimates produced by the ACF method, as non-stationary data can lead to inconsistent results in
400 autocorrelation-based analyses.

3.4.3 Contribution of instrumental noise in the measured LOS velocity variances

~~This section evaluates, beam by beam, the impact of instrumental noise on the measured LOS velocity variances, which are~~
~~combined to derive the along-wind variance (Eqs. 5-6).~~ The parametric expression (Eq. 2) used to fit the LOS velocity spectra
measured by beam i enables the identification of the power spectral density of instrumental noise, N_i , and the derivation of
the variances, $\sigma_{n_i}^2$ (Eq. 4). Fig. 9 compares the mean magnitude of $\sigma_{n_i}^2$ to the mean variance of the net contribution from
405

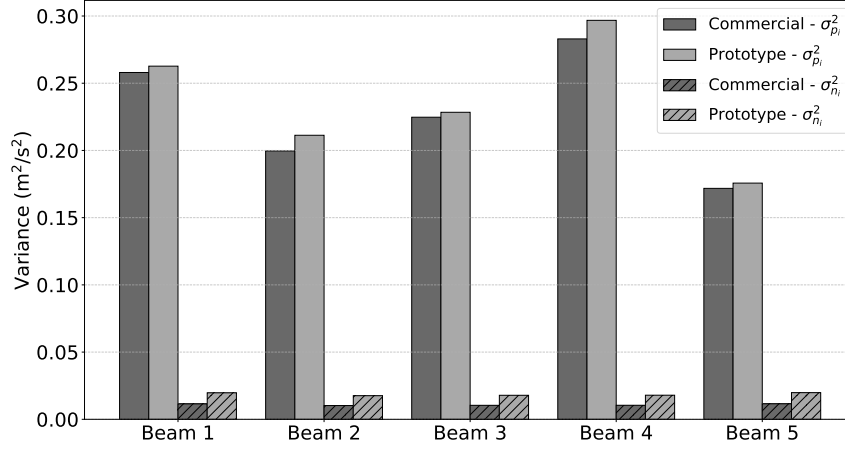


Figure 9. Mean variance of the net contribution from atmospheric turbulence ($\sigma_{p_i}^2$), corrected for instrumental noise derived from the spectral method, measured by each beam i at scales observable by the commercial (dark gray) and prototype (light gray) lidar profilers. Dashed areas represent the mean variance of instrumental noise, $\sigma_{n_i}^2$. The averages were computed over the 47-day dataset.

atmospheric turbulence, $\sigma_{p_i}^2$, corrected for instrumental noise at scales observable by the commercial and prototype lidar profilers.

The mean values of $\sigma_{n_i}^2$, which are nearly identical across all beams, were found to be $0.0108 \text{ m}^2/\text{s}^2$ for the commercial configuration (Table 4). A similar trend was observed for the prototype configuration, although the mean variance of instrumental noise was 68% higher, at $0.0181 \text{ m}^2/\text{s}^2$ (Table 4). Notably, the contribution of instrumental noise variance to the total variance, $\sigma_{b_i}^2$ (Eq. 3), was found to be 4.8% and 7.4% for the commercial and prototype lidar configurations, respectively.

The mean variances, $\sigma_{p_i}^2$ were consistently higher for measurements obtained with the prototype configuration. Across all beams, the mean value was $0.2288 \text{ m}^2/\text{s}^2$, which is 7.8% higher than the corresponding mean value for the commercial lidar measurements.

Table 4. Median and mean (\pm spread) variance of instrumental noise for commercial and prototype lidars, computed from the LOS velocity measurements across all beams using spectral and ACF methods.

	Commercial lidar		Prototype lidar	
Methods	Spectral	ACF	Spectral	ACF
Median (m^2/s^2)	0.0076	0.0050	0.0129	0.0081
Mean \pm spread (m^2/s^2)	0.0108 ± 0.0102	0.0148 ± 0.0228	0.0181 ± 0.0175	0.0237 ± 0.0294

Table 5. Error Statistics of the along-wind standard deviation derived from the variance method, corrected for instrumental noise, applied on measurements collected by the commercial and prototype lidars in comparison to the reference sonic anemometer.

	Bias (m/s)	MAE (m/s)	RMSE (m/s)	R ²	Relative Error (%)
Commercial lidar	-0.0639	0.0886	0.1218	0.9138	7.8
Prototype lidar	-0.0466	0.0678	0.0871	0.9574	5.7

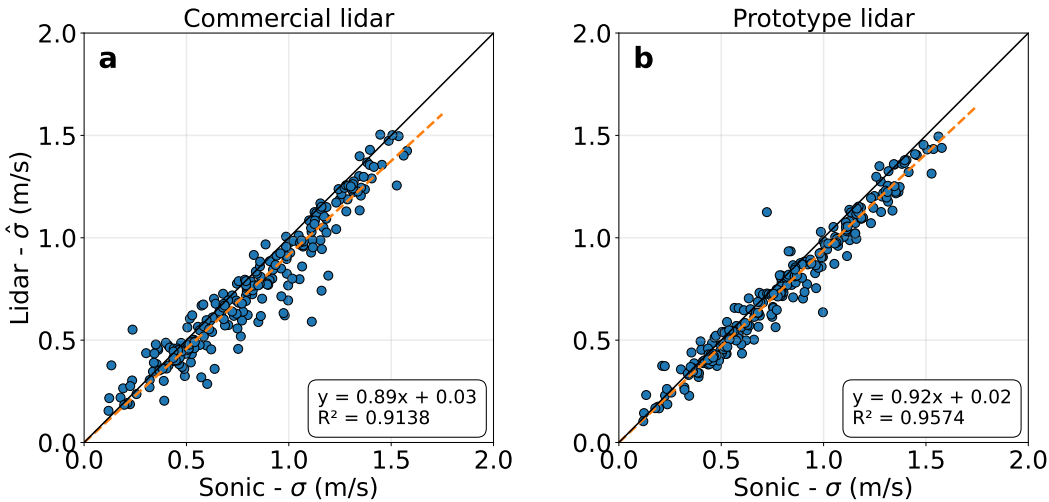


Figure 10. Scatter plots of along-wind standard deviation, $\hat{\sigma}$, derived from the variance method applied on measurements of the commercial and prototype lidar configurations versus standard deviation, σ , derived from the reference sonic anemometer. The standard deviation estimates are restricted to cases where wind direction was aligned with one pair of opposite beams.

415 **3.5 Along-wind standard deviation**

Fig. 10 presents scatter plots of the along-wind standard deviation, $\hat{\sigma}$, derived from the variance method applied on measurements of both lidar configurations compared to the standard deviation, σ , obtained from the reference sonic anemometer. The prototype configuration demonstrates superior performance across all error metrics, with bias, MAE, and RMSE approximately 25% lower than those of the commercial configuration (Table 5). Additionally, the coefficient of determination is 5% higher. There is also a reduction in the relative error of the mean standard deviation, with the prototype configuration showing values of 5.7% compared to 7.8% for the commercial configuration.

420

Fig. 11 presents bin-averaged-wind speed-binned estimates of $\hat{\sigma}$ compared to estimates of σ (black curve) as a function of binned-averaged wind speed. For all wind speeds, the standard deviation measured by the sonic anemometer consistently remains higher than that derived from both lidar configurations. Below wind speed of 8 m/s, the standard deviation values from both lidar configurations closely match each other. Within this wind speed range, the standard deviation associated with the commercial lidar is 2.7% higher than that from the prototype configuration. However, above this wind speed threshold,

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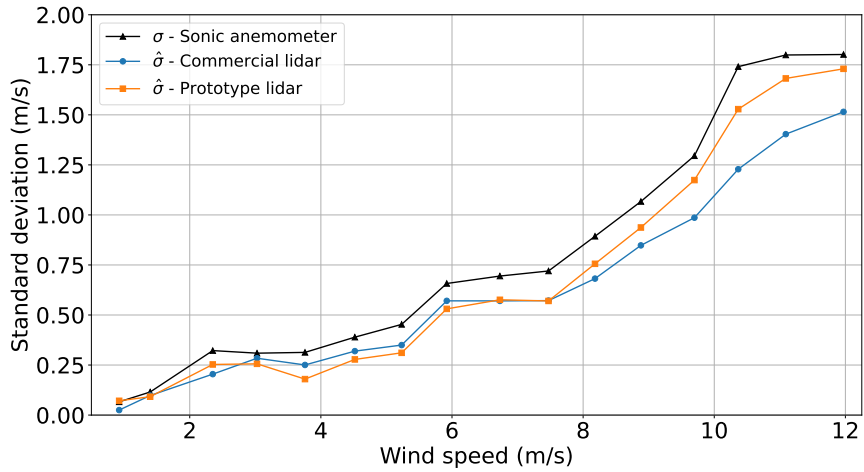


Figure 11. Along-wind standard deviation, $\hat{\sigma}$, derived from measurements of the commercial lidar (blue curve) and prototype lidar (orange curve), compared to the standard deviation, σ , obtained from reference sonic anemometer measurements (black curve) as a function of wind speed.

the standard deviation associated with the prototype configuration increases more rapidly with wind speed compared to the commercial lidar. In this wind speed range, the standard deviation associated with the prototype lidar is 13.0% higher than that associated with the commercial configuration. For all wind speed ranges, the prototype lidar measurements exhibited a mean standard deviation and variance that were 2.9% and 3.5% and 7.2% higher, respectively, than those of the commercial configuration.

~~Scatter plots of along-wind standard deviation, $\hat{\sigma}$, derived from the variance method applied on measurements of the commercial and prototype lidar configurations versus standard deviation, σ , derived from the reference sonic anemometer. The standard deviation estimates are restricted to cases where wind direction was aligned with one pair of opposite beams.~~

4 Discussion

~~When proposing an increase in sampling rate to capture smaller eddies and their associated energy through variance, it is essential to assess the instrumental noise and its variance to ensure that the observed changes are due to physical phenomena rather than noise. In this study, we estimate noise-induced variance using two distinct methods. For both lidar configurations, the mean variances of instrumental noise computed from two methods were found to be consistent with values obtained in previous studies, such as the WindCube lidar analysis by Mann et al. (2009). This alignment reinforces confidence in our estimates.~~

The increased sampling rate leads to higher instrumental noise compared to the commercial configuration, as expected, since the noise variance is inversely proportional to the number of transmitted pulses (Pearson et al., 2009). In the prototype lidar, achieving a higher sampling rate required reducing the number of pulses leading to the elevated noise levels. The noise variance

445 was approximately 5% of the total variance for the commercial configuration and over 7% for the prototype configuration. While the noise contribution is relatively low, it is not negligible, and its impact should be considered when calculating second-order statistics of LOS velocities in pulsed wind lidar profilers.

Compared to estimates derived from the commercial lidar configuration, the error metrics (bias, MAE, RMSE) of the along-wind standard deviation estimates corrected for instrumental noise using the prototype lidar were notably lower. Additionally, 450 the mean along-wind variance measured by the prototype was higher, suggesting that the increased sampling rate allows for improved detection of turbulent energy associated with smaller eddies. However, this observed improvement remains significantly below the theoretical benefit expected from increasing the LOS sampling rate, as determined through sonic anemometer measurements. The measurement volume of a sonic anemometer is effectively point-like, in comparison to the much larger probe length of wind lidar profilers. The anemometer is in fact essentially free of the intra-beam effect, which 455 enables it to capture the wind signature of very small eddies.

One way to mitigate the impact of the intra-beam effect is by significantly reducing the probe length. A shorter probe length would minimize spatial averaging, enhancing the resolution of high-frequency fluctuations in the wind field. However, this reduction is likely to weaken the lidar echo strength, potentially increasing measurement errors. Therefore, a balance must be struck between improving resolution and maintaining signal quality.

460 The increased sampling rate of the prototype configuration could positively impact metrics such as turbulence intensity (TI), which is a key factor due to its influence on structural loads and turbine lifespan. As demonstrated in Fig. 11, this effect is expected to be more pronounced at higher wind speeds, such as at 15 m/s. At this wind speed, TI—referred to as TI15—is used in design turbulence calculations to define characteristic standard deviation bands, which are essential for classifying turbines according to wind turbulence, as outlined in IEC 61400-1.

465 Moreover, the increased sampling rate extends the LOS velocity spectra over the frequency domain. This is particularly valuable for floating lidar systems, where measurements are used to derive TI through motion-compensation algorithms which is an ongoing research topic (e.g., Kelberlau et al., 2020; Désert et al., 2021). The buoy's motion generates energy spikes in the LOS velocity spectra at frequencies corresponding to the wave periods, the buoy's natural period, and their interactions (Thiébaud et al., 2024a). This introduces additional variance, $\sigma_{m_i}^2$, which must be accounted for alongside the three other 470 variance terms that contribute to the total variance, $\sigma_{b_i}^2$, measured by each beam i (Eq. 3). With a Nyquist frequency of $f_{N_c} = 0.125$ Hz, the commercial configuration can detect motion-induced variance for periods up to 8 seconds, while lower-period motion remains undetectable. In contrast, the prototype configuration presented in this paper, with a Nyquist frequency of $f_{N_p} = 0.5$ Hz, enables the detection of motion effects for periods as short as 2 seconds. Thus, the prototype configuration may offer an advantage in developing motion-compensation algorithms based on the variance method by improving the derivation of true 475 variance through LOS velocity spectra analysis.

The increased sampling rate resulted in a relatively slight 0.5% reduction in data availability compared to the commercial configuration over the 47-day dataset. While this difference is minimal, it may become more noticeable over longer measurement campaigns, which typically last over a year for wind site characterization. Following the measurement campaign presented in this paper, the prototype configuration was installed in December 2022 on Planier Island in the Mediterranean Sea, where

480 it remains operational. The wind characteristics derived from the full year of 2023 are presented in Thiébaud et al. (2024b), including a detailed analysis of data availability. Encouragingly, up to 160 m above sea level, annual data availability exceeded the 90% threshold considered best practice. Beyond this height, availability gradually declined, reaching below 70% at 220 m. While this highlights an area for further optimization, the prototype lidar has already demonstrated strong performance at critical measurement heights.

485 Moreover, the prototype configuration performed comparably to the commercial setup in terms of mean wind characteristics. While the commercial configuration met the ~~"best-practice"~~ "best practice" threshold for all key performance indicators (KPIs), the prototype also achieved this standard, with the exception of mean wind speed differences, where it met the ~~"minimum acceptance"~~ "minimum acceptance" level within the best practice range. This result is promising, as it confirms that the prototype lidar meets industry standards while offering opportunities for further refinement. ~~With continued development,~~
490 ~~the prototype lidar has the potential to further enhance wind resource assessments and support the needs of modern wind energy projects.~~

5 Conclusions

~~This study highlights both the potential and challenges of enhancing the~~

Reducing the accumulation time increases the sampling rate, which helps limit temporal averaging of the wind signal and
495 preserves more of the high-frequency variance within the instrument's resolvable range. However, this also requires a careful evaluation of instrumental noise and its associated variance to ensure that observed changes in variance can be confidently attributed to atmospheric turbulence rather than measurement artifacts. In this study, the noise-induced variance was estimated using two independent methods. For both lidar configurations, the noise levels were consistent with values reported in previous studies, such as the WindCube lidar analysis by Mann et al. (2009), supporting the reliability of our estimates. As expected,
500 ~~the higher~~ the higher ~~sampling rate in wind lidar systems, particularly for turbulence measurements and the detection of energy from smaller eddies. The prototype lidar configuration, with its increased sampling rate, demonstrated advantages in capturing higher-frequency fluctuations in the wind field, leading to more accurate along-wind variance estimates. By resolving smaller eddies, the prototype effectively captured additional turbulent energy that the commercial configuration, with its lower sampling rate, could not detect. This improvement was especially pronounced~~ the prototype lidar led to increased instrumental noise due
505 to the reduced number of transmitted pulses (Pearson et al., 2009). The noise variance constituted approximately 5% of the total variance for the commercial configuration and over 7% for the prototype. While non-negligible, this contribution was accounted for in all variance-based metrics.

However, the impact of increased sampling must be interpreted with care. If the sampling frequency is too low relative to the turbulent fluctuations present in the flow, aliasing occurs: unresolved high-frequency energy is folded into lower frequency
510 bands, distorting the spectral distribution. While the reduced accumulation time increases the sampling frequency and helps mitigate aliasing effects, it does not fully recover the true spectral shape, given the finite temporal resolution. Therefore, the

observed increase in variance with the prototype lidar should primarily be attributed to reduced temporal and spatial filtering, rather than a direct "gain" in turbulent energy capture. This distinction is critical.

515 Increasing the sampling rate does not linearly "capture more variance" from smaller eddies. Instead, the accumulation time acts as a low-pass filter on the LOS velocity signal, attenuating contributions from high-frequency turbulent fluctuations. Systems with longer accumulation times, such as the 0.8 s used in the commercial lidar, are more affected by this filtering, especially at higher wind speeds, ~~where the prototype's reduced minimum detectable eddy size provided significant benefits. However, the theoretical gains expected from where the increased sampling rate were not fully realized due to inherent measurement limitations within the probe, where spatial and temporal filtering effects constrain turbulence detection at certain scales.~~

520 ~~The~~ effective probe length, L_{eff} , becomes longer due to advection. The effective probe length refers to the spatial distance over which the LOS wind velocity is effectively averaged, accounting for both the fixed probe length and the distance traveled by air during the accumulation period. It can be estimated as:

$$L_{\text{eff}} = L_{\text{probe}} + U\Delta t, \quad (11)$$

525 A longer effective probe length increases spatial averaging and attenuates variability in the measurements. In contrast, the prototype lidar, with a shorter 0.2 s accumulation time, experiences less temporal averaging and thus preserves a greater portion of the wind variance from smaller-scale motions. This contributes to the higher along-wind variances observed in the prototype configuration compared to the commercial system. Despite this, the observed increase in variance remains significantly below the theoretical benefit expected from increasing the LOS sampling rate, as determined through sonic anemometer measurements. The measurement volume of a sonic anemometer is effectively point-like, in comparison to the

530 much larger probe length of wind lidar profilers. The anemometer is essentially free of the probe-time averaging effect, which enables it to capture the wind signature of very small eddies.

5 Conclusions

The prototype configuration of the WindCube v2.1 lidar profiler demonstrated comparable performance to the commercial system in terms of mean wind characteristics and data availability, meeting industry standards. Moreover, this study demonstrates

535 that increasing the sampling rate of LOS wind velocity measurements by reducing the accumulation time in pulsed wind lidar systems effectively mitigates temporal and spatial averaging effects inside the probe volume, thereby improving turbulence measurements. The prototype configuration, with a fourfold increase in sampling rate compared to the commercial system, preserved additional variance by reducing both spatial and temporal averaging, rather than directly resolving smaller eddies, particularly at higher wind speeds, where advection lengthens the effective probe length. Despite these gains, the increased

540 sampling rate ~~also introduced some introduced~~ trade-offs, including ~~elevated higher~~ instrumental noise and a slight reduction in data availability. ~~Noise contributions to the total variances were non-negligible and required correction to ensure accurate turbulence statistics. It is worth noting that a configuration similar to the prototype examined in this study could be achieved~~

with a commercial lidar, as manufacturers can program an increased sampling rate with relative ease. However, users should be aware that implementing such a modification would require validation against a meteorological mast to obtain certification, which could add time and cost to deployment. The noise was systematically corrected, and its impact was found to be manageable.

Another key consideration is balancing increased sampling rate with reduced probe length. Ideally, combining both enhancements would yield the optimal configuration, improving both temporal and spatial resolution. In this study, the increased sampling rate proved beneficial for turbulence measurements while maintaining acceptable accuracy for mean wind statistics. Meanwhile, the effects of reducing the probe length are currently under investigation. While further research is needed to determine its feasibility in field applications, this approach holds promise for future advancements in wind lidar profiler performance. The findings underscore the importance of carefully balancing temporal resolution, noise, and probe length when configuring lidar systems for turbulence retrieval. While commercial lidars can be programmed by the manufacturer to match the prototype's sampling rate, it is advisable to validate the lidar's performance through certification after such adjustments. Looking forward, optimizing both accumulation time and probe length in tandem may enhance the ability of wind lidar systems to capture turbulence more accurately. However, such changes must be approached cautiously, as reducing these parameters can also compromise data availability, especially under challenging atmospheric conditions. Balancing resolution and reliability will be key to supporting broader applications in wind energy assessments and atmospheric boundary layer research.

Author contributions

MT identified the problematic, performed the analysis and drafted the paper. LM performed the mathematical demonstration of Eq. 1, available in the supplementary material, and provided guidance on the manuscript structure and its review. FD and FG reviewed the manuscript.

Data and code availability

The data is owned by a private consortium with proprietary rights and confidentiality obligations, precluding its sharing alongside this paper.

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570 **Competiting interest**

The authors declare that they have no conflict of interest.

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References

- Bodini, N., Lundquist, J. K., and Newsom, R. K.: Estimation of turbulence dissipation rate and its variability from sonic anemometer and wind Doppler lidar during the XPIA field campaign, *Atmospheric Measurement Techniques*, 11, 4291–4308, 2018.
- Bodini, N., Lundquist, J. K., Krishnamurthy, R., Pekour, M., Berg, L. K., and Choukulkar, A.: Spatial and temporal variability of turbulence
580 dissipation rate in complex terrain, *Atmospheric Chemistry and Physics*, 19, 4367–4382, 2019.
- Branlard, E., Pedersen, A. T., Mann, J., Angelou, N., Fischer, A., Mikkelsen, T., Harris, M., Slinger, C., and Montes, B. F.: Retrieving wind statistics from average spectrum of continuous-wave lidar, *Atmospheric Measurement Techniques*, 6, 1673–1683, 2013.
- Browning, K. A. and Wexler, R.: The determination of kinematic properties of a wind field using Doppler radar, *Journal of Applied meteorology and climatology*, 7, 105–113, 1968.
- 585 Brugger, P., Träumner, K., and Jung, C.: Evaluation of a procedure to correct spatial averaging in turbulence statistics from a Doppler lidar by comparing time series with an ultrasonic anemometer, *Journal of Atmospheric and Oceanic Technology*, 33, 2135–2144, 2016.
- DNV and NORSEWInD: Best Practice Test and Verification Procedure for Wind Lidars on the Høvsøre Test Site, Tech. Rep. GL GH-D Report WT 6960/09 for EU-Project NORSEWInD, 2009.
- Doviak, R. J. and Zrníc, D. S.: *Doppler Radar & Weather Observations*, Courier Corporation, Courier Corporation, 1993.
- 590 Durgesh, V., Thomson, J., Richmond, M. C., and Polagye, B. L.: Noise correction of turbulent spectra obtained from acoustic doppler velocimeters, *Flow Measurement and Instrumentation*, 37, 29–41, 2014.
- Désert, T., Knapp, G., and Aubrun, S.: Quantification and correction of wave-induced turbulence intensity bias for a floating lidar system, *Remote Sensing*, 13, 2973, 2021.
- Eberhard, W. L., Cupp, R. E., and Healy, K. R.: Doppler lidar measurement of profiles of turbulence and momentum flux, *Journal of*
595 *Atmospheric and Oceanic Technology*, 6, 809–819, 1989.
- Frehlich, R.: Coherent Doppler lidar signal covariance including wind shear and wind turbulence, *Applied optics*, 33, 6472–6481, publisher: Optical Society of America, 1994.
- Gal-Chen, T., Xu, M., and Eberhard, W. L.: Estimations of atmospheric boundary layer fluxes and other turbulence parameters from Doppler lidar data, *Journal of Geophysical Research: Atmospheres*, 97, 18 409–18 423, 1992.
- 600 Gargett, A. E., Tejada-Martinez, A. E., and Grosch, C. E.: Measuring turbulent large-eddy structures with an ADCP. Part 2. Horizontal velocity variance, 2009.
- Guerra, M. and Thomson, J.: Turbulence Measurements from Five-Beam Acoustic Doppler Current Profilers, *Journal of Atmospheric and Oceanic Technology*, 34, 1267–1284, 2017.
- Hayat, M. J.: *Understanding statistical significance*, Nursing research, 59, 219–223, publisher: LWW, 2010.
- 605 Kelberlau, F. and Mann, J.: Cross-contamination effect on turbulence spectra from Doppler beam swinging wind lidar, *Wind Energy Science*, 5, 519–541, 2020.
- Kelberlau, F., Neshaug, V., Lønseth, L., Bracchi, T., and Mann, J.: Taking the motion out of floating lidar: Turbulence intensity estimates with a continuous-wave wind lidar, *Remote Sensing*, 12, 898, 2020.
- Kristensen, L., Kirkegaard, P., and Mikkelsen, T.: Determining the velocity fine structure by a laser anemometer with fixed orientation,
610 *Danmarks Tekniske Universitet, Risø Nationallaboratoriet for Bæredygtig Energi*, 2011.
- Kropfli, R. A.: Single Doppler radar measurements of turbulence profiles in the convective boundary layer, *Journal of Atmospheric and Oceanic Technology*, 3, 305–314, 1986.

- Lenschow, D. H., Mann, J., and Kristensen, L.: How long is long enough when measuring fluxes and other turbulence statistics?, *Journal of Atmospheric and Oceanic Technology*, 11, 661–673, 1994.
- 615 Lenschow, D. H., Wulfmeyer, V., and Senff, C.: Measuring second-through fourth-order moments in noisy data, *Journal of Atmospheric and Oceanic technology*, 17, 1330–1347, 2000.
- Lhermitte, R. M.: Note on wind variability with Doppler radar, *Journal of Atmospheric Sciences*, 19, 343–346, 1962.
- Lhermitte, R. M.: Note on the observation of small-scale atmospheric turbulence by Doppler radar techniques, *Radio Science*, 4, 1241–1246, 1969.
- 620 Lu, Y. and Lueck, R. G.: Using a broadband ADCP in a tidal channel. Part II: Turbulence, *Journal of Atmospheric and Oceanic Technology*, 16, 1568–1579, 1999.
- Mann, J., Cariou, J.-P., Courtney, M. S., Parmentier, R., Mikkelsen, T., Wagner, R., Lindelow, P., Sjöholm, M., and Enevoldsen, K.: Comparison of 3D turbulence measurements using three staring wind lidars and a sonic anemometer, *Meteorologische Zeitschrift*, 18, 135, 2009.
- 625 Mann, J., Peña, A., Bingöl, F., Wagner, R., and Courtney, M. S.: Lidar scanning of momentum flux in and above the atmospheric surface layer, *Journal of Atmospheric and Oceanic Technology*, 27, 959–976, 2010.
- McMillan, J. M. and Hay, A. E.: Spectral and structure function estimates of turbulence dissipation rates in a high-flow tidal channel using broadband ADCPs, *Journal of Atmospheric and Oceanic Technology*, 34, 5–20, 2017.
- Newman, J. F., Klein, P. M., Wharton, S., Sathe, A., Bonin, T. A., Chilson, P. B., and Muschinski, A.: Evaluation of three lidar scanning
630 strategies for turbulence measurements, *Atmospheric Measurement Techniques*, 9, 1993–2013, 2016.
- Olesen, H. R., Larsen, S. E., and Højstrup, J.: Modelling velocity spectra in the lower part of the planetary boundary layer, *Boundary-Layer Meteorology*, 29, 285–312, 1984.
- O'Connor, E. J., Illingworth, A. J., Brooks, I. M., Westbrook, C. D., Hogan, R. J., Davies, F., and Brooks, B. J.: A method for estimating the turbulent kinetic energy dissipation rate from a vertically pointing Doppler lidar, and independent evaluation from balloon-borne in situ
635 measurements, *Journal of atmospheric and oceanic technology*, 27, 1652–1664, 2010.
- Pearson, G., Davies, F., and Collier, C.: An analysis of the performance of the UFAM pulsed Doppler lidar for observing the boundary layer, *Journal of Atmospheric and Oceanic Technology*, 26, 240–250, 2009.
- Peña, A., Hasager, C. B., Gryning, S., Courtney, M., Antoniou, I., and Mikkelsen, T.: Offshore wind profiling using light detection and ranging measurements, *Wind Energy*, 12, 105–124, 2009.
- 640 Richard, J.-B., Thomson, J., Polagye, B., and Bard, J.: Method for identification of doppler noise levels in turbulent flow measurements dedicated to tidal energy, *International Journal of Marine Energy*, 3, 52–64, 2013.
- Rippeth, T. P., Williams, E., and Simpson, J. H.: Reynolds stress and turbulent energy production in a tidal channel, *Journal of Physical Oceanography*, 32, 1242–1251, 2002.
- Sathe, A. and Mann, J.: Measurement of turbulence spectra using scanning pulsed wind lidars, *Journal of Geophysical Research: Atmospheres*, 117, 2012.
645
- Sathe, A., Mann, J., Gottschall, J., and Courtney, M. S.: Can wind lidars measure turbulence?, *Journal of Atmospheric and Oceanic Technology*, 28, 853–868, 2011.
- Sjöholm, M., Mikkelsen, T., Mann, J., Enevoldsen, K., and Courtney, M.: Spatial averaging-effects on turbulence measured by a continuous-wave coherent lidar, *Meteorologische Zeitschrift (Berlin)*, 18, 2009.

650 Smalikho, I., Köpp, F., and Rahm, S.: Measurement of atmospheric turbulence by 2- μ m Doppler lidar, *Journal of Atmospheric and Oceanic Technology*, 22, 1733–1747, 2005.

Stacey, M. T., Monismith, S. G., and Burau, J. R.: Measurements of Reynolds stress profiles in unstratified tidal flow, *Journal of Geophysical Research*, 104, 10 935–10 949, 1999a.

Stacey, M. T., Monismith, S. G., and Burau, J. R.: Observations of turbulence in a partially stratified estuary, *Journal of Physical*
655 *Oceanography*, 29, 1950–1970, publisher: American Meteorological Society, 1999b.

Strauch, R. G., Merritt, D. A., Moran, K. P., Earnshaw, K. B., and De Kamp, D. V.: The Colorado wind-profiling network, *Journal of Atmospheric and Oceanic Technology*, 1, 37–49, 1984.

Stull, R. B.: *Meteorology for scientists and engineers: a technical companion book with Ahrens' Meteorology Today*, 2000.

Teunissen, H. W.: Structure of mean winds and turbulence in the planetary boundary layer over rural terrain, *Boundary-Layer Meteorology*,
660 19, 187–221, 1980.

Theriault, K.: Incoherent multibeam Doppler current profiler performance: Part II–Spatial response, *IEEE journal of oceanic engineering*, 11, 16–25, 1986.

Thiébaud, M., Filipot, J.-F., Maisondieu, C., Damblans, G., Duarte, R., Droniou, E., Chaplain, N., and Guillou, S.: A comprehensive assessment of turbulence at a tidal-stream energy site influenced by wind-generated ocean waves, *Energy*, 191, 116 550, 2020.

665 Thiébaud, M., Quillien, N., Maison, A., Gaborieau, H., Ruiz, N., MacKenzie, S., Connor, G., and Filipot, J.-F.: Investigating the flow dynamics and turbulence at a tidal-stream energy site in a highly energetic estuary, *Renewable Energy*, 195, 252–262, 2022.

Thiébaud, M., Thebault, N., Le Boulluec, M., Damblans, G., Maisondieu, C., Benzo, C., and Guinot, F.: Experimental Evaluation of the Motion-Induced Effects for Turbulent Fluctuations Measurement on Floating Lidar Systems, *Remote Sensing*, 16, 1337, 2024a.

Thiébaud, M., Vonta, L., Benzo, C., and Guinot, F.: Characterization of the offshore wind dynamics for wind energy production in the Gulf
670 of Lion, Western Mediterranean Sea, *Wind Energy and Engineering Research*, 1, 100 002, 2024b.

Thomson, J., Polagye, B., Durgesh, V., and Richmond, M. C.: Measurements of turbulence at two tidal energy sites in Puget Sound, WA, *Oceanic Engineering*, *IEEE Journal of Oceanic Engineering*, 37, 363–374, 2012.

Tieleman, H. W.: Universality of velocity spectra, *Journal of Wind Engineering and Industrial Aerodynamics*, 56, 55–69, 1995.

Welch, P.: The use of fast Fourier transform for the estimation of power spectra: A method based on time averaging over short, modified
675 periodograms, *IEEE Transactions on audio and electroacoustics*, 15, 70–73, 1967.

Wildmann, N., Bodini, N., Lundquist, J. K., Bariteau, L., and Wagner, J.: Estimation of turbulence dissipation rate from Doppler wind lidars and in situ instrumentation for the Perdigão 2017 campaign, *Atmospheric Measurement Techniques*, 12, 6401–6423, 2019.

Wilson, D. A.: *Doppler radar studies of boundary layer wind profile and turbulence in snow conditions*, 1970.